

Molecular Cloud Fragmentation: Self-Gravity, Magnetic Fields, Ambipolar Diffusion, and Nonlinear Flows

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Star Formation Rate

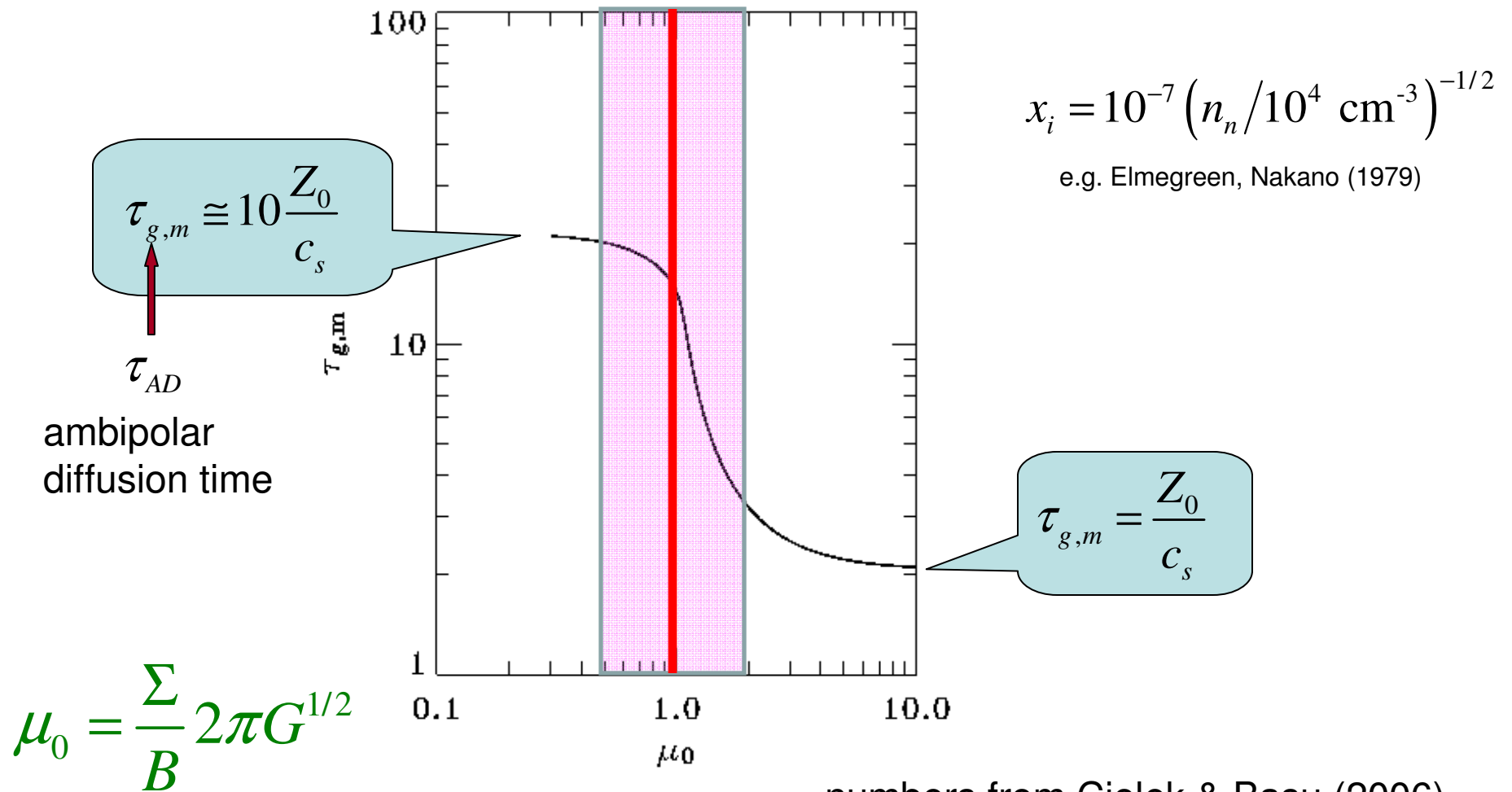
Regulated (to what extent) by:

$$\Delta \mathbf{v}, B, x_i$$

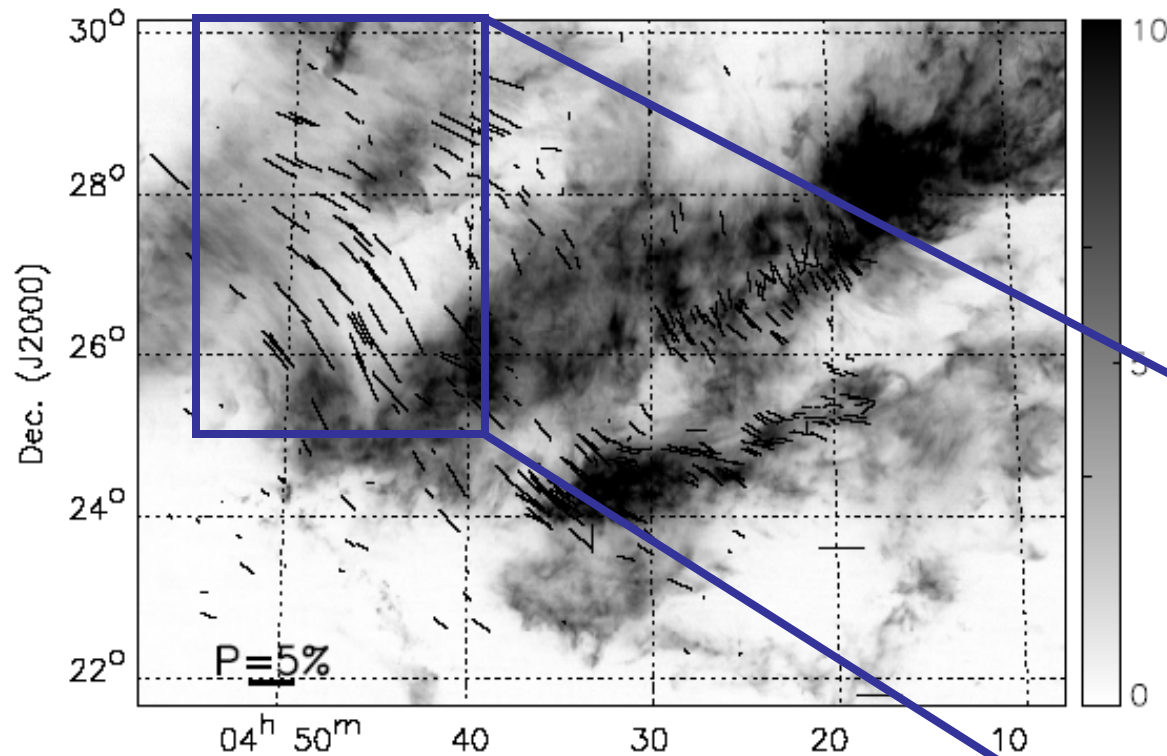
i.e., nonlinear/nonthermal motions (turbulence, MHD waves)
magnetic field strength (really mass-to-flux ratio $\mu = 2\pi G^{1/2} \Sigma / B$)
ionization fraction (sets neutral-ion coupling)

$\mu \approx 1$ (transcritical) is interesting

For CR ionized sheet, with half thickness Z_0 .



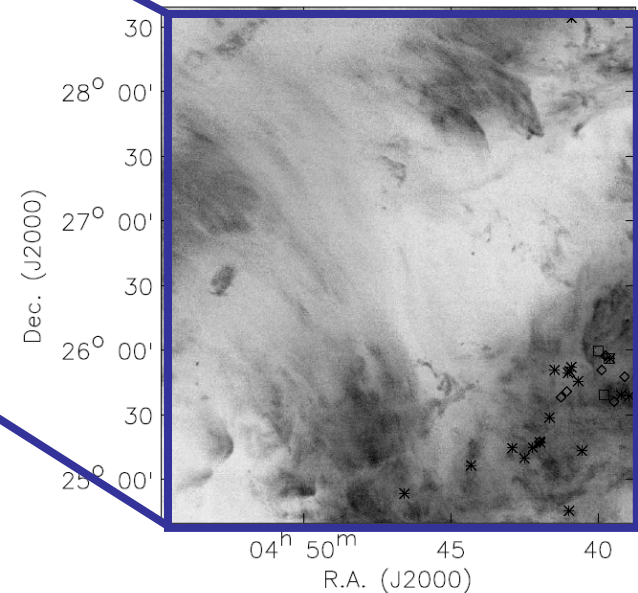
Taurus – *B* dominated envelope?



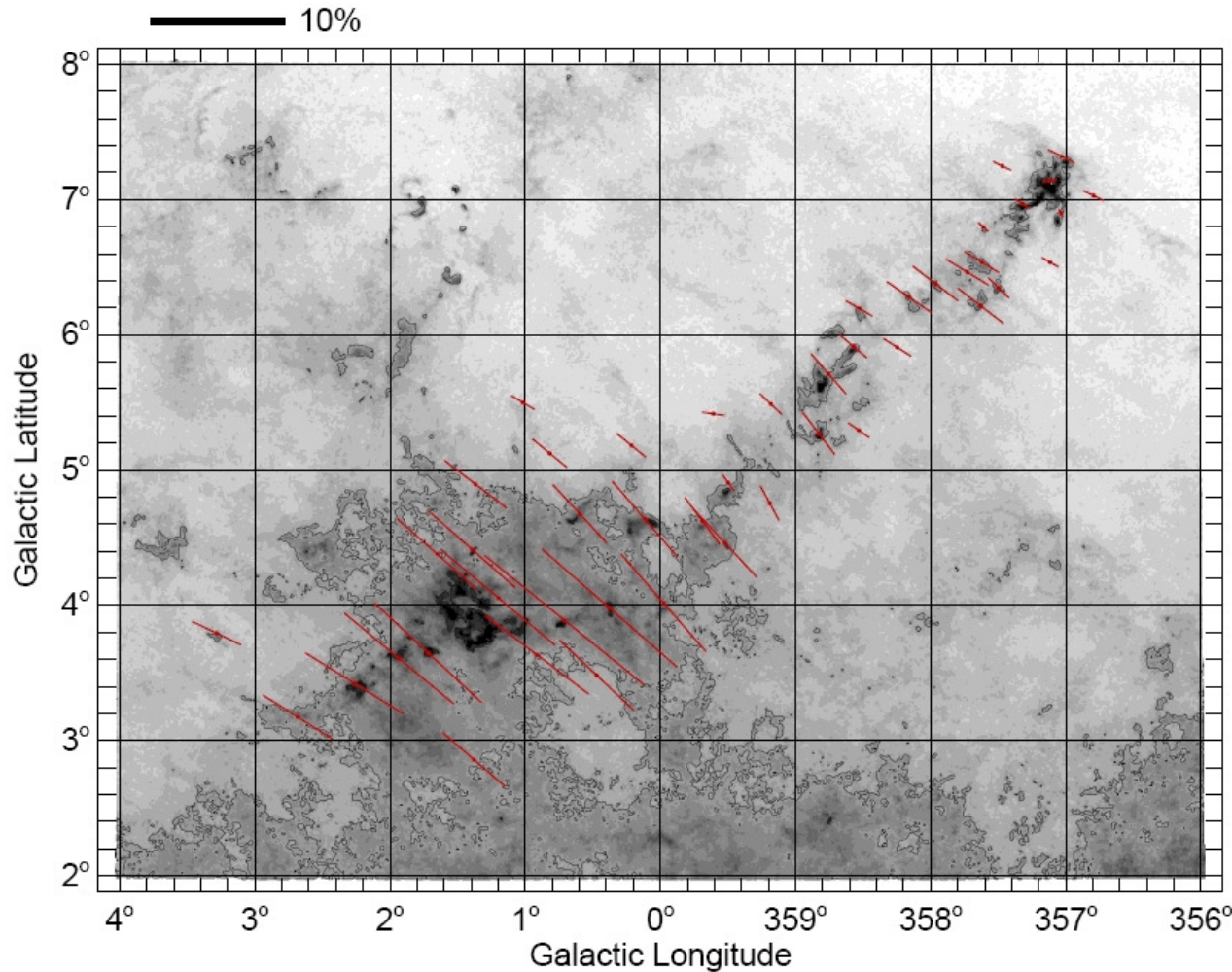
Goldsmith et al. (2008), ^{12}CO emission

Striations of gas emission consistent with magnetically-dominated envelope.

Heyer et al. (2008): Pol. maps + vel. info \rightarrow low plasma beta in envelope \Leftrightarrow subcritical? Most mass is in low density envelope (Goldsmith et al. 08), so probably, yes.



Pipe Nebula – more B effect?

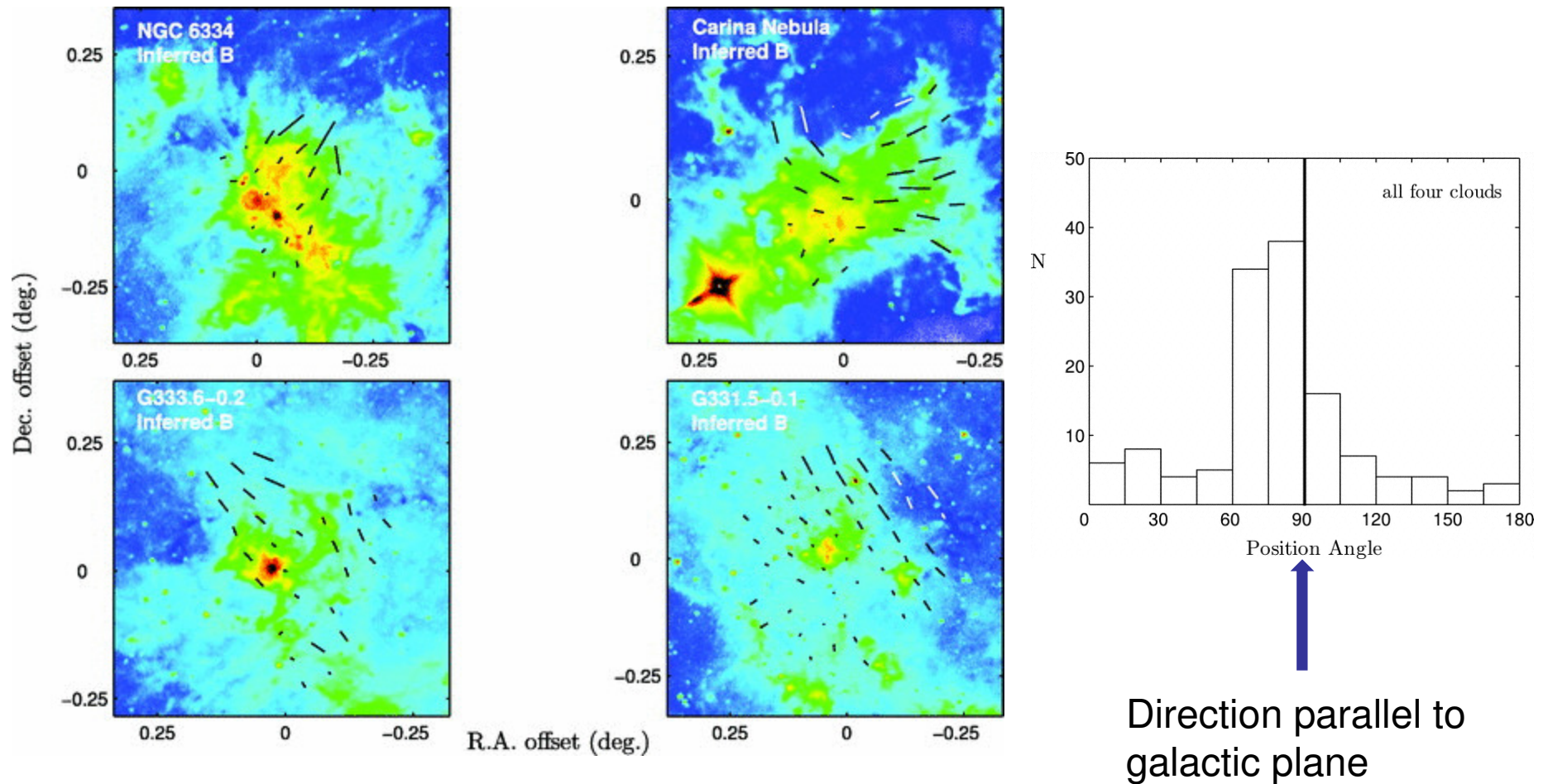


Magnetically
regulated cloud
formation?

Pipe (and Taurus)
→ formed by flow or
contraction along B ?

Alves, Franco, &
Girart (2008)

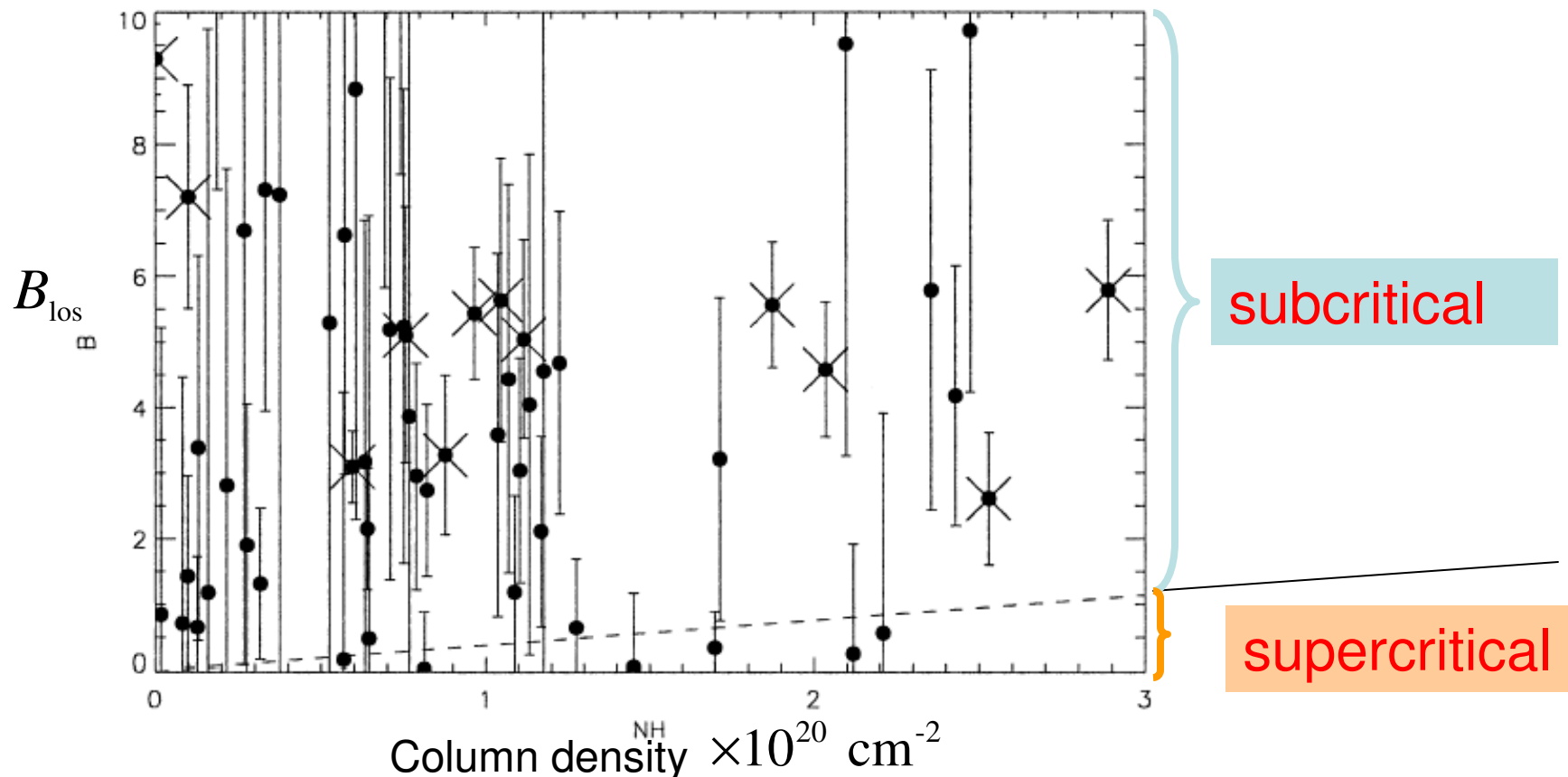
GMC Fields align with Galactic B



H. Li et al. (2006)

MC progenitors are H I Clouds

Heiles & Troland (2008)



Flux freezing in H I gas → Critical or supercritical MC formation requires significant accumulation of mass ALONG the magnetic field.

Hard to accumulate supercritical MC rapidly

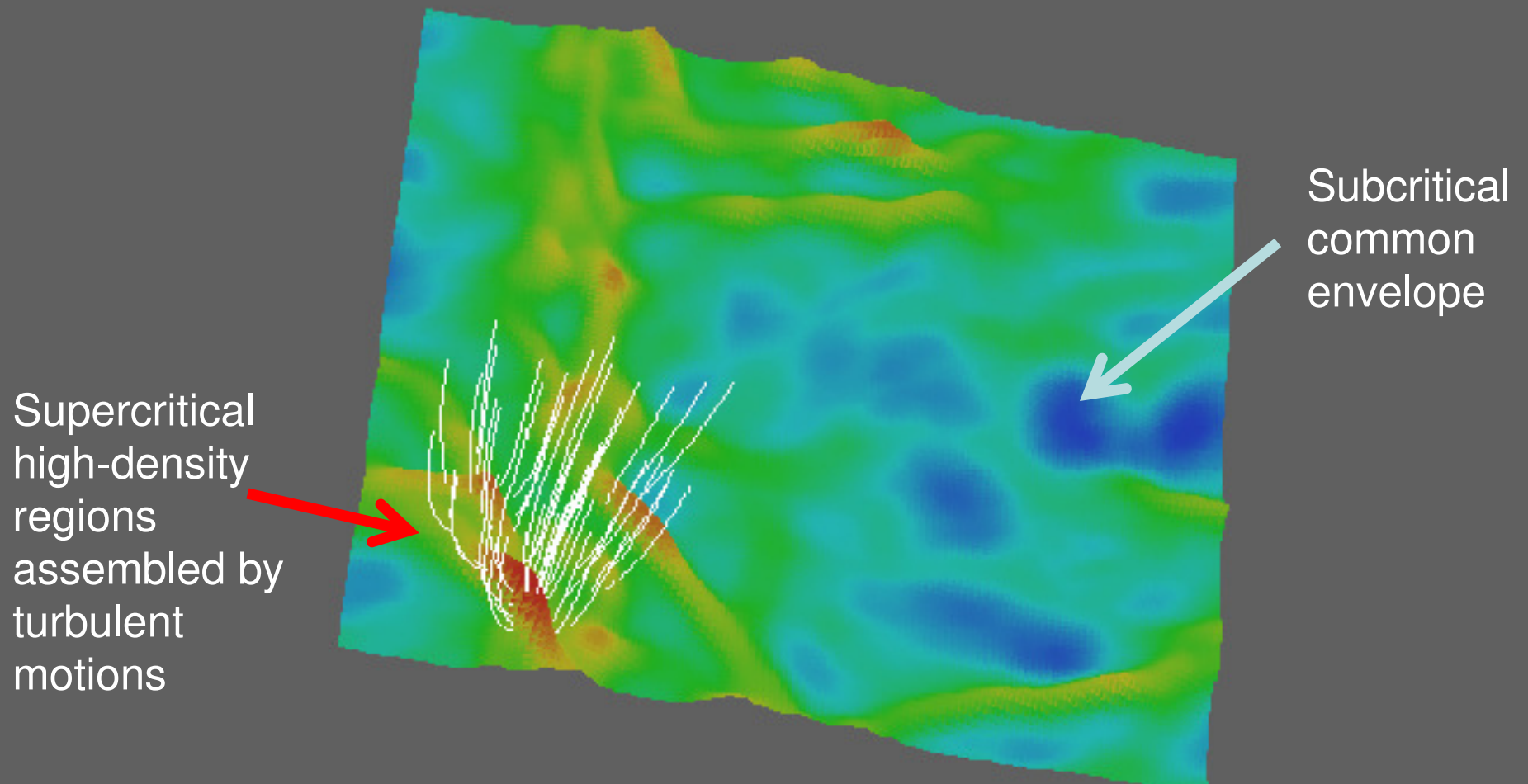
$$L \cong \frac{B}{2\pi G^{1/2} \rho} \cong 150 \left(\frac{\mu}{1} \right) \left(\frac{B}{3 \times 10^{-6} \text{ G}} \right) \left(\frac{n}{1 \text{ cm}^{-3}} \right) \text{ pc},$$

$$v \equiv \frac{L}{t} \cong 150 \left(\frac{\mu}{1} \right) \left(\frac{B}{3 \times 10^{-6} \text{ G}} \right) \left(\frac{n}{1 \text{ cm}^{-3}} \right) \left(\frac{t}{10^6 \text{ yr}} \right)^{-1} \text{ km/s}.$$

Mestel (1999), Stellar Magnetism, and earlier papers quote even larger value, 10^3 above, not 150.

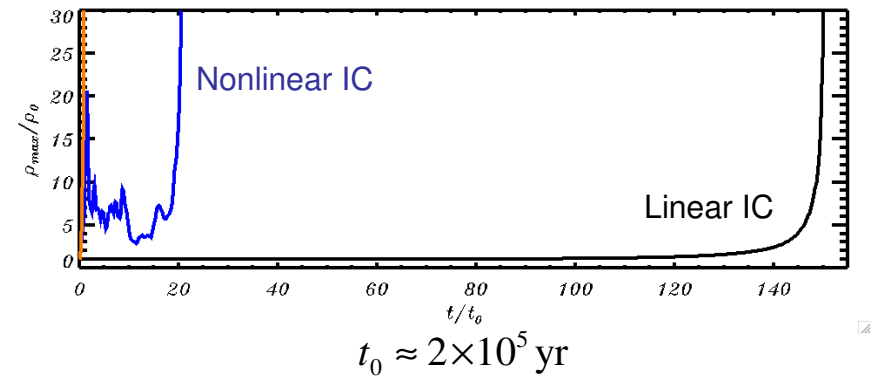
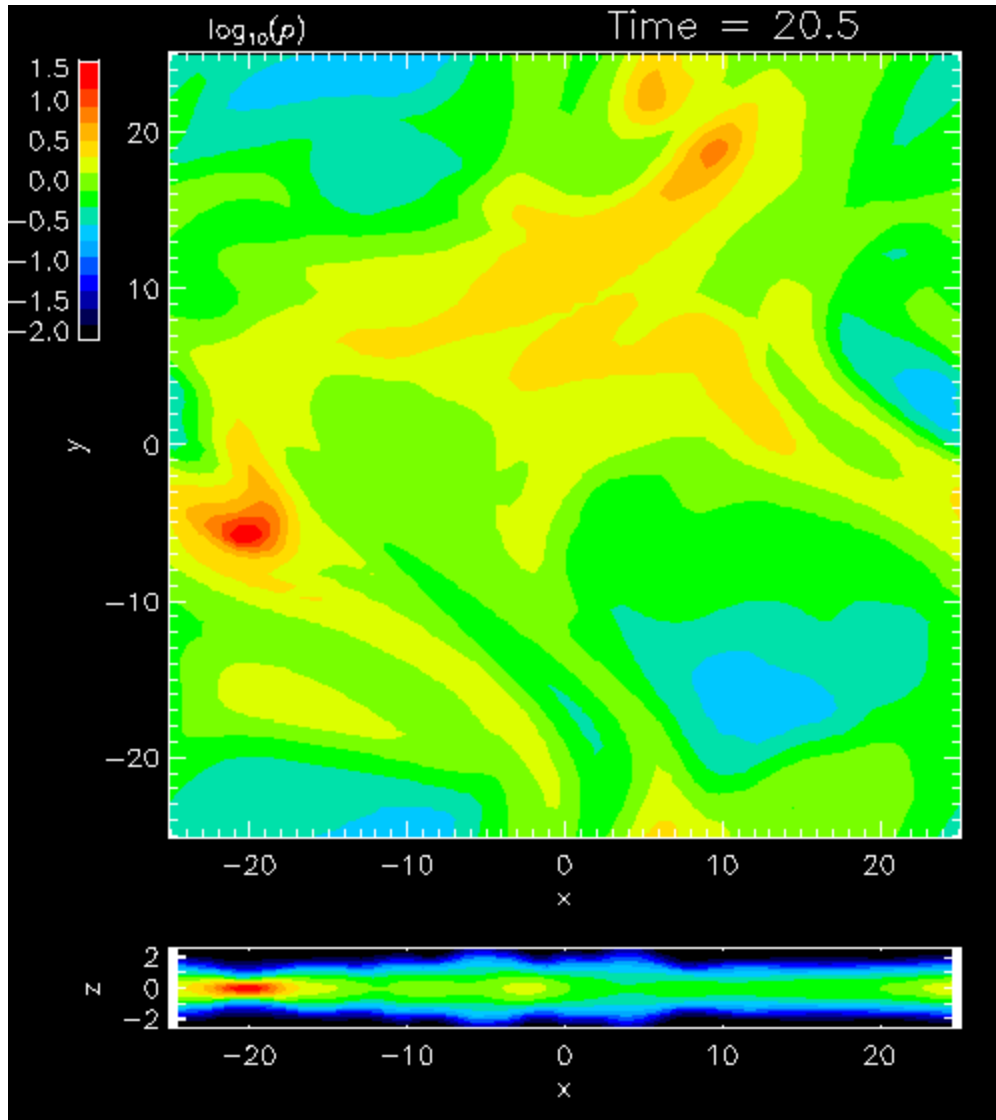
Bottom line: Highly supercritical MC **AND** rapid formation time t is troublesome!

Scenario



cf. Nakamura & Li (2005), Elmegreen (2007), Kudoh & Basu (2008), Nakamura & Li (2008), Basu, Ciolek, Dapp, & Wurster (2009; model shown above).

3D Trans-Alfvénic Model → Rapid but Inefficient SF



Nonlinear initial velocity field

$$v_k^2 \propto k^{-4}$$

allowed to decay

rms amplitude $v_a = 3c_s \approx v_{\text{Alf}}$

trans-Alfvénic

$$\mu_0 \equiv 2\pi G^{1/2} \Sigma_0 / B_0 = 0.5$$

$$\text{ionization } x_i = 10^{-7} (n_n / 10^4 \text{ cm}^{-3})^{-1/2}$$

Gas density in midplane ($z=0$)

A vertical slice of gas density

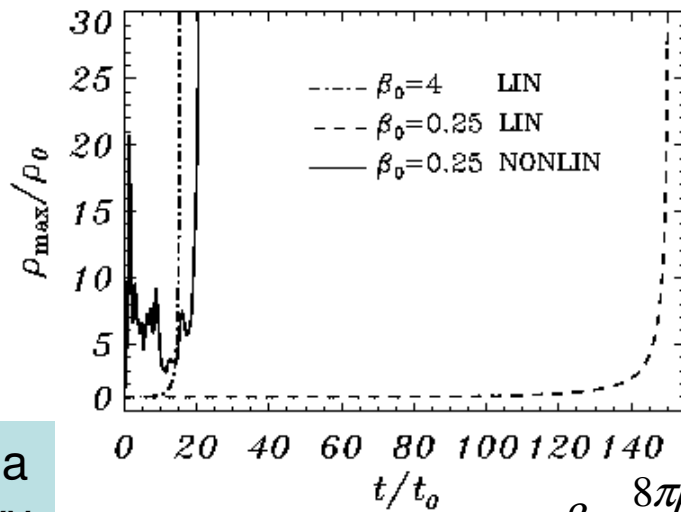
box width = 2.5 pc

Originally 64 x 64 x 40 cells, now 256 x 256 x 40

Kudoh & Basu (2008, ApJ, 679, L79)

How Does Turbulent Ambipolar Diffusion Work?

Runaway collapse of the first core occurs ~ 7 times faster with nonlinear as opposed to small-amplitude IC's.



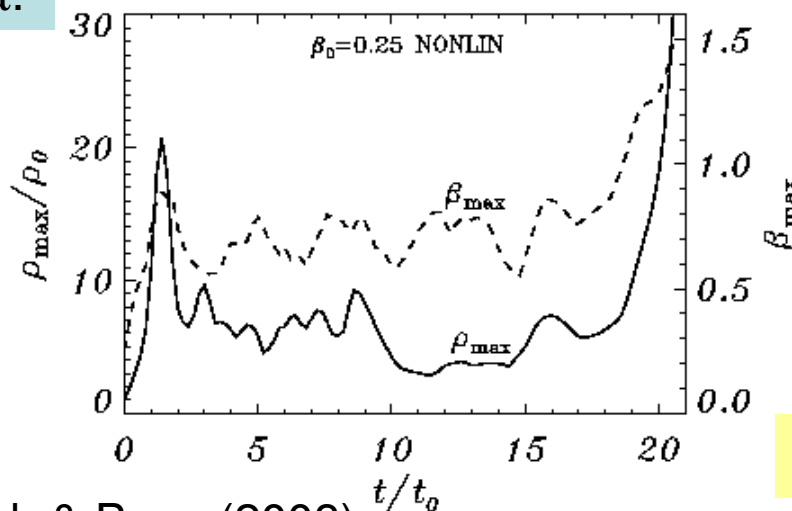
Early turbulent compression

$$\tau_{AD} \propto L^{5/2} \Rightarrow \beta \uparrow \text{ quickly as } L \downarrow$$

β is a proxy for μ .

$$\beta \equiv \frac{8\pi\rho c_s^2}{B} \propto \mu^2$$

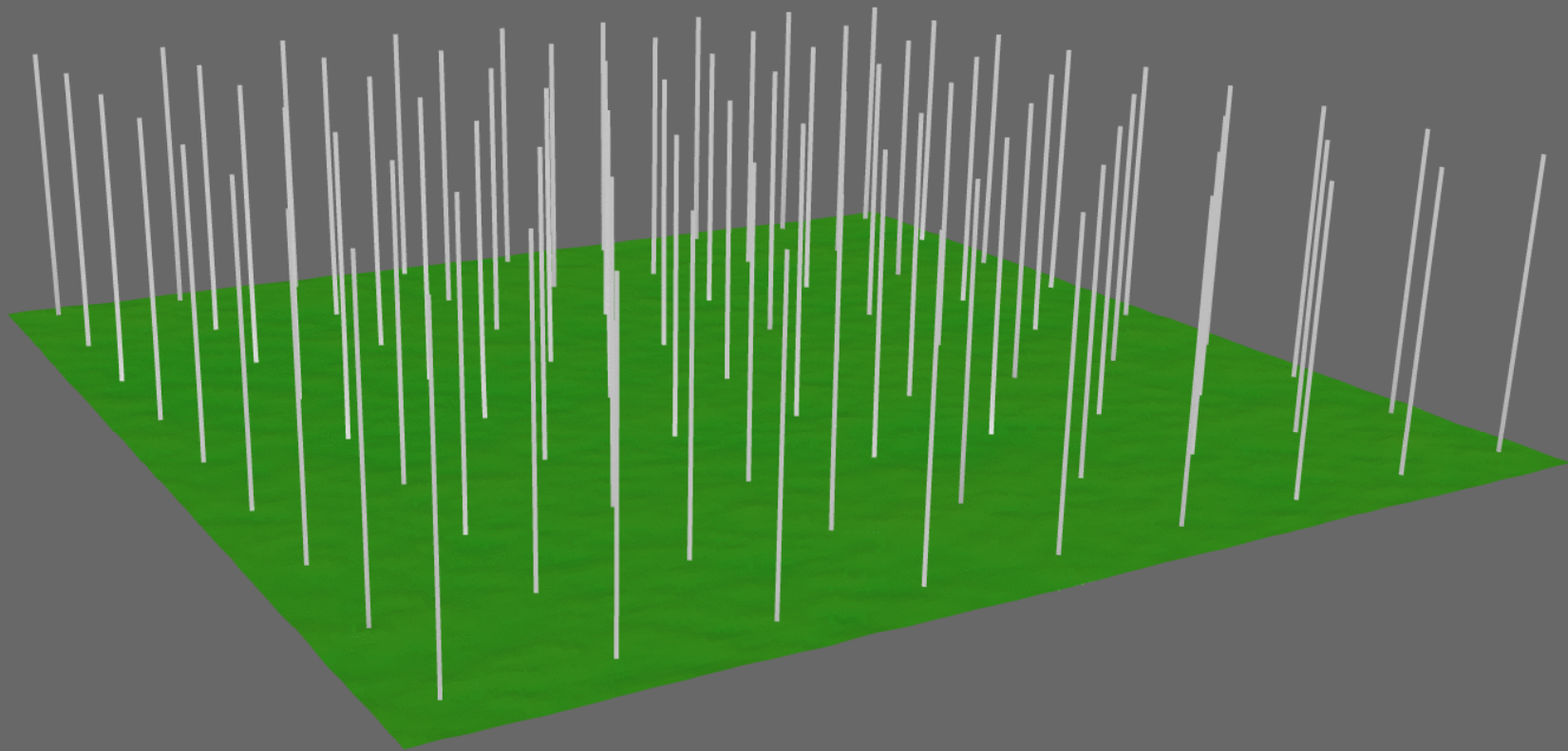
Later evolution at higher mean density than in initial state



$\beta \uparrow$ continues more slowly

Rapid contraction when/where $\beta > 1$.

Animation with Field Lines



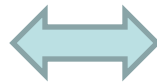
Thin disk approx. – from models of Basu, Ciolek, Dapp, & Wurster (2009, NewA, 14, 483)

SFR Related to Timescale for Runaway Collapse

Accelerated collapse
(some oscillations
before collapse)

$$t_{coll} \geq 2\pi Z_0 / c_s$$

\approx system dynamical time



$$v_a < v_{MS,0}$$

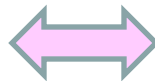
initial rms
turbulent speed

initial magnetosonic
speed

Prompt collapse

$$t_{coll} \approx Z_0 / v_a$$

less than dynamical time

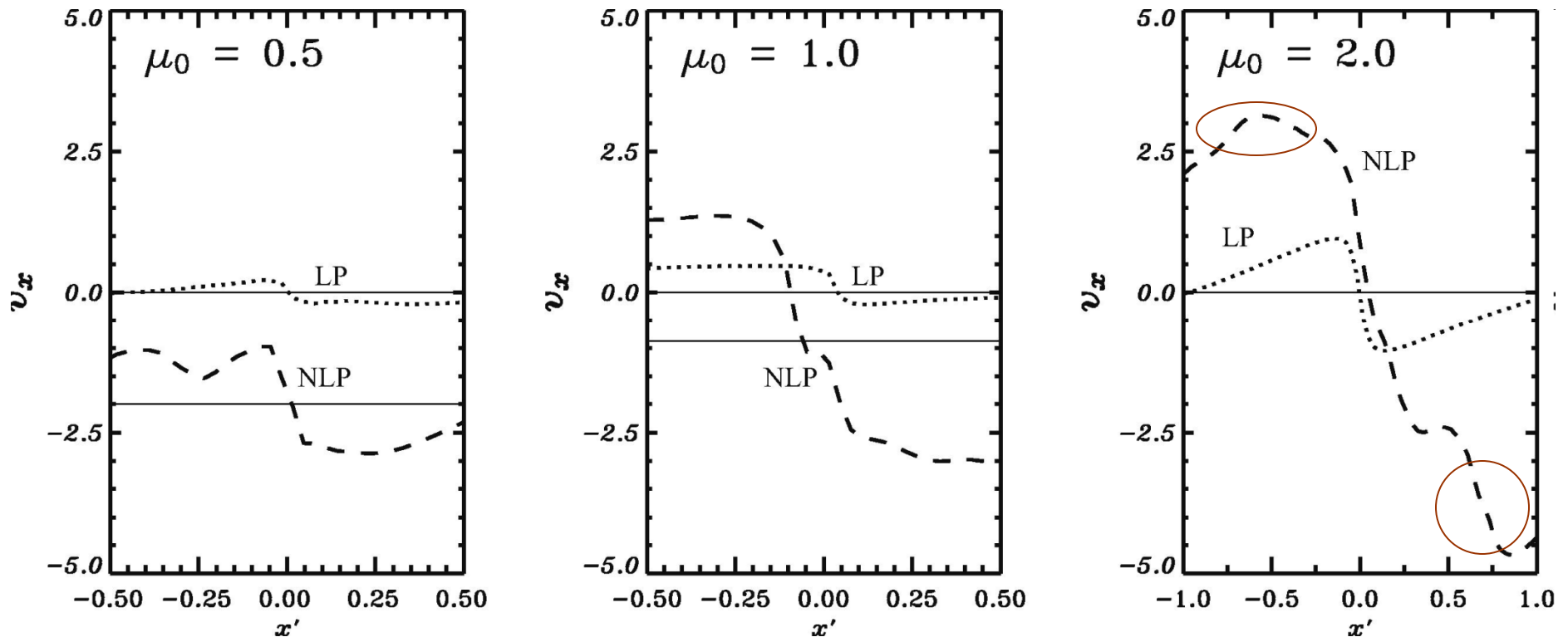


$$v_a > v_{MS,0}$$

for ionization fraction $x_i = 10^{-7} (n_n / 10^4 \text{ cm}^{-3})^{-1/2}$

Expect much greater SFR, especially
for driven super-Alfvénic turbulence.

Super-Alfvénic Model Fails



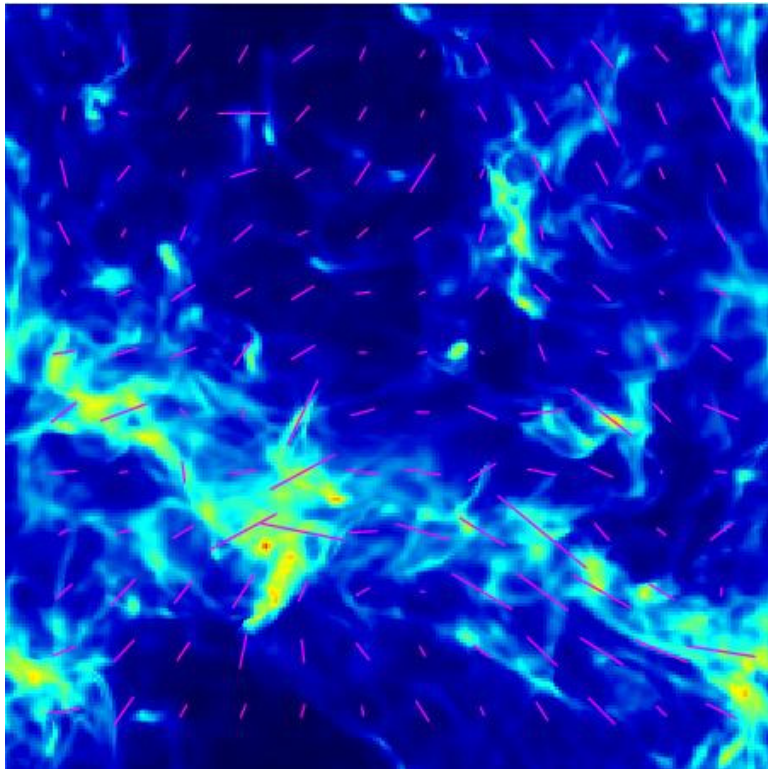
Velocity cuts through cores

LP = initial Linear Perturbations

NLP = initial Nonlinear Perturbations $v_a = 2c_s \rightarrow$ Super-Alfvénic for $\mu_0 = 2$.

Super Alfvénic \rightarrow highly supersonic infall and immediate SF.

Turbulent Decay

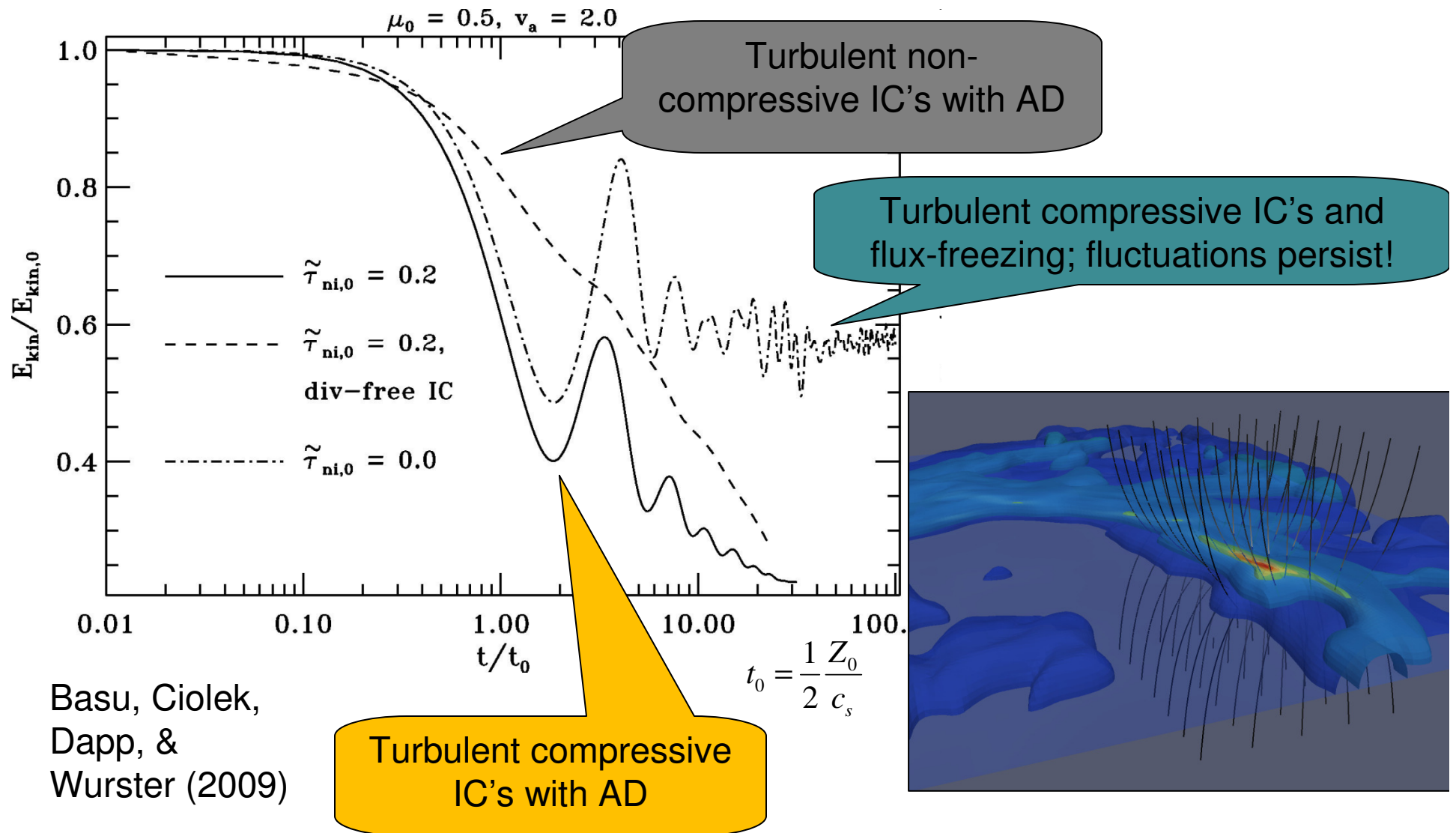


3D local (periodic box) simulations.
Left: integrated column density and simulated polarization maps (Ostriker et al. 1999)

Important result: turbulence decays away on crossing time of characteristic length scale. A robust result for infinite uniform medium.

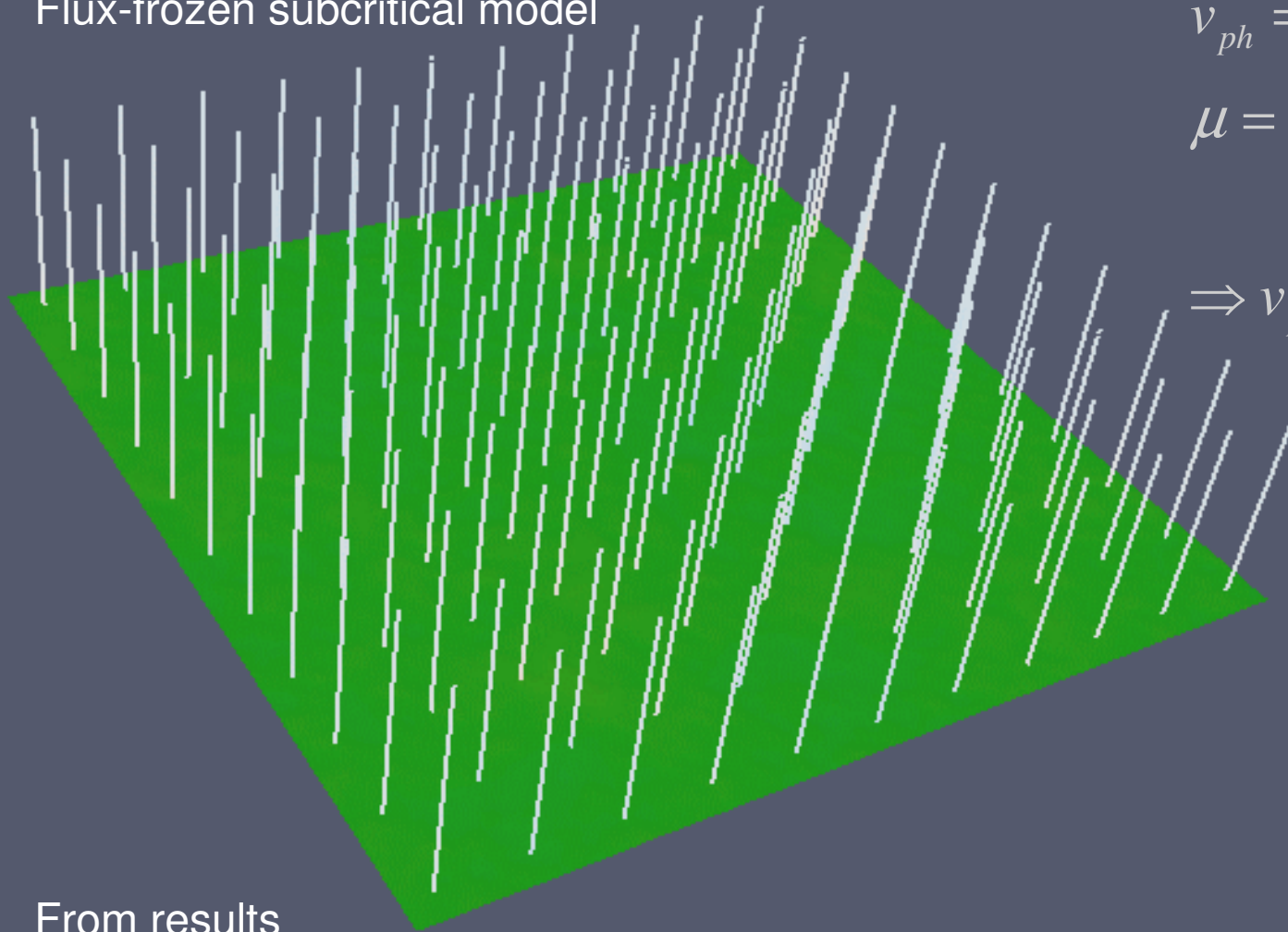
Stone, Ostriker, & Gammie (1998), MacLow et al. (1998), Ostriker, Stone, and Gammie (1999; image above) and many others.

Turbulent Decay in our models – usually fast, but...



How can fluctuations persist?

Flux-frozen subcritical model



$$v_{ph} = \sqrt{(\mu^{-2} - 1)G \Sigma \lambda},$$

$$\mu = 0.5, \lambda = 8c_s^2 / G \Sigma$$

(box size)

$$\Rightarrow v_{ph} \approx 5c_s.$$

From results
of BCDW09

Results and Emerging Scenario

- Thin disk model yields long lived supersonic motions. Effectively flux-frozen (**UV ionized**, e.g. McKee 1989) cloud envelopes can, **if subcritical**, maintain indefinite oscillations due to restoring force of externally-anchored B . Future...fully global models.
- Meanwhile, 3D models \rightarrow inner CR ionized regions can undergo turbulence accelerated ambipolar diffusion \rightarrow rapid core formation.
- SF Rate implications: Trans-Alfvénic turbulence leads to accelerated but low efficiency SF. Super-Alfvénic turbulence leads to prompt SF with very high velocity flows in core vicinity – ruled out generally.

More: see poster by Wolf Dapp