

SF@50

5 -10 July, 2009, Siena, Italy

Relations between the Global SFR and Intrinsic Structures of the ISM in Galactic Disks

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- 1) Euler grid modeling galactic disks
- 2) SPH modeling galactic disks
- 3) Euler grid modeling a molecular disk around an SMBH

Q1: What determines the global SFR?

A₁: Gravitational, thermal, magneto-rotational or hydrodynamic instabilities in a gas disk

A₂: Turbulence in molecular clouds

A₃: Collisions between molecular clouds

A₄: Molecular fraction in dense ISM

A₅: Galactic rotation

A₆: Galactic spiral

A₇: H₂ chemistry, FUV

....

Q2: Then, which is the most important?

A: **All.** So we need a 5-day workshop!

Today's talk = A simple theoretical picture on global SFR

Combination of these processes produces a mess in a galactic disk ==> a robust statistical feature in the ISM ==> KS-like SF law

Numerical Experiment of global evolution of galactic disks using a hydrodynamic code:

HD.AUSM.MUSCL.Cartesian.Uniform.
Convolution.FFT.radiative_cool.Zsun
.photoelectric_heating.SN.neat.spher
ical_pot.SMBH.Xray.Conductio
n.H2

e.g.

tracking stellar particle: HD.*.STAR

H2 formation/dissociation: HD.*.H2

Spiral potential: HD.*.spiral_pot

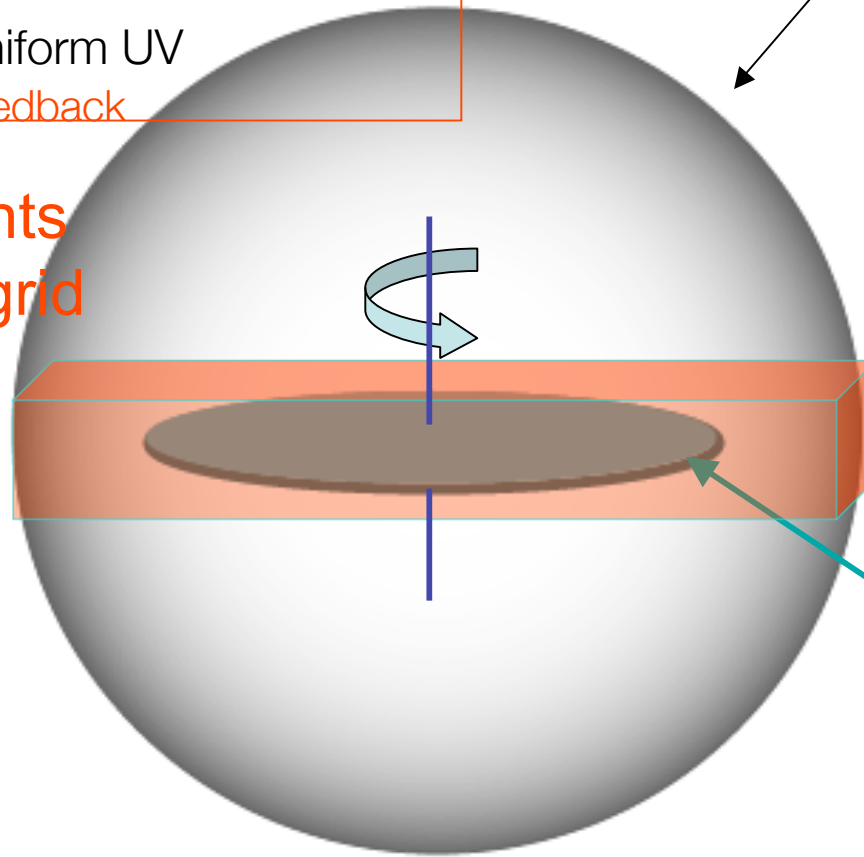
Xray heating: HD.*.Xray

3-D Hydrodynamics of a gas disk in a spherical galactic potential

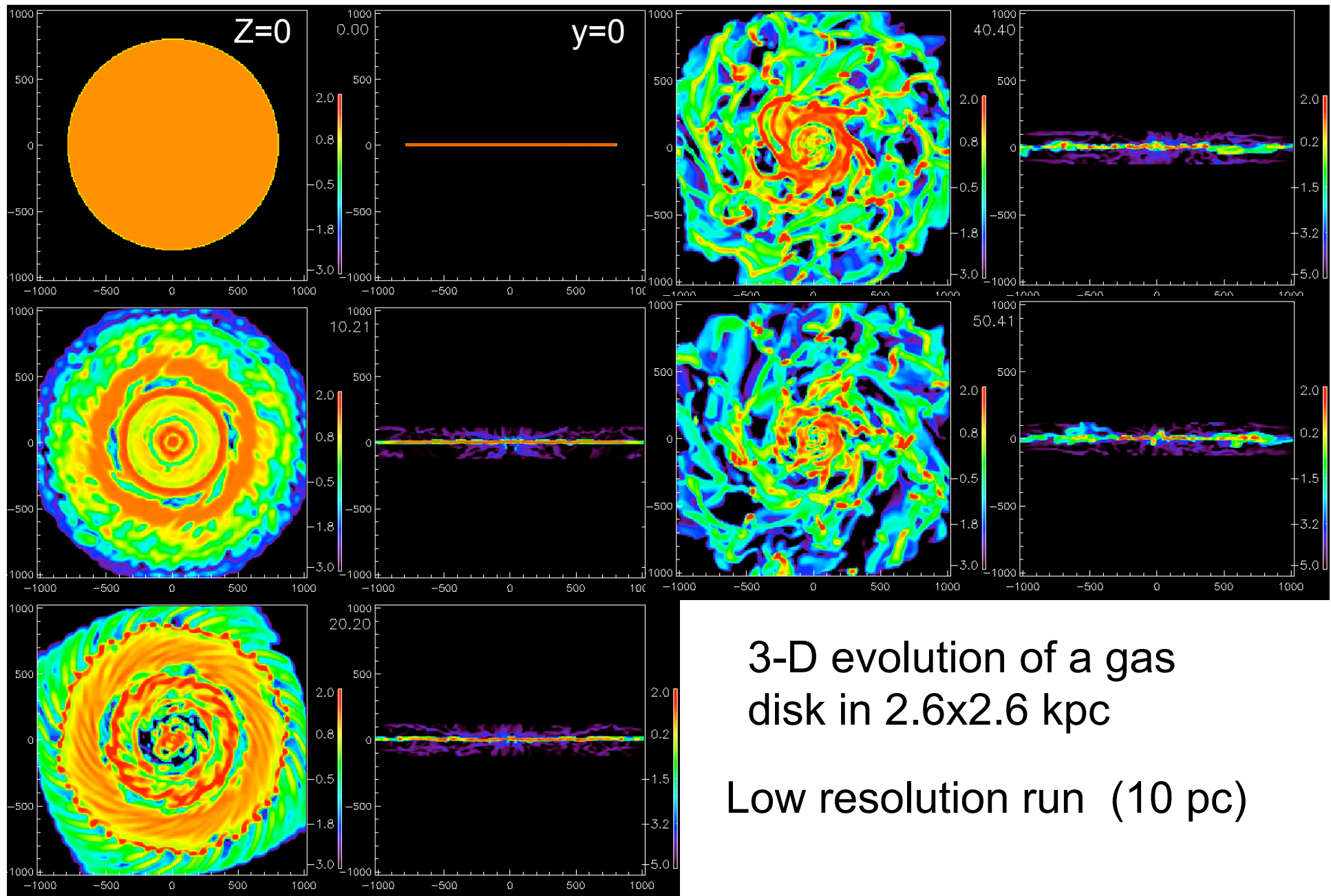
- Self-gravity of the gas
- A cooling function ($10 < T < 10^8$ K) is assumed (solar)
- Heating sources: uniform UV
 - No SNe energy feedback

Stellar/DM potential: fixed

512²x64 grid points
Resolution: 5pc/grid
(also 10pc/grid)



Thin gas disk
 $\sim 10^{8-9} M_{\text{sun}}$
 $R = 1.28$ kpc

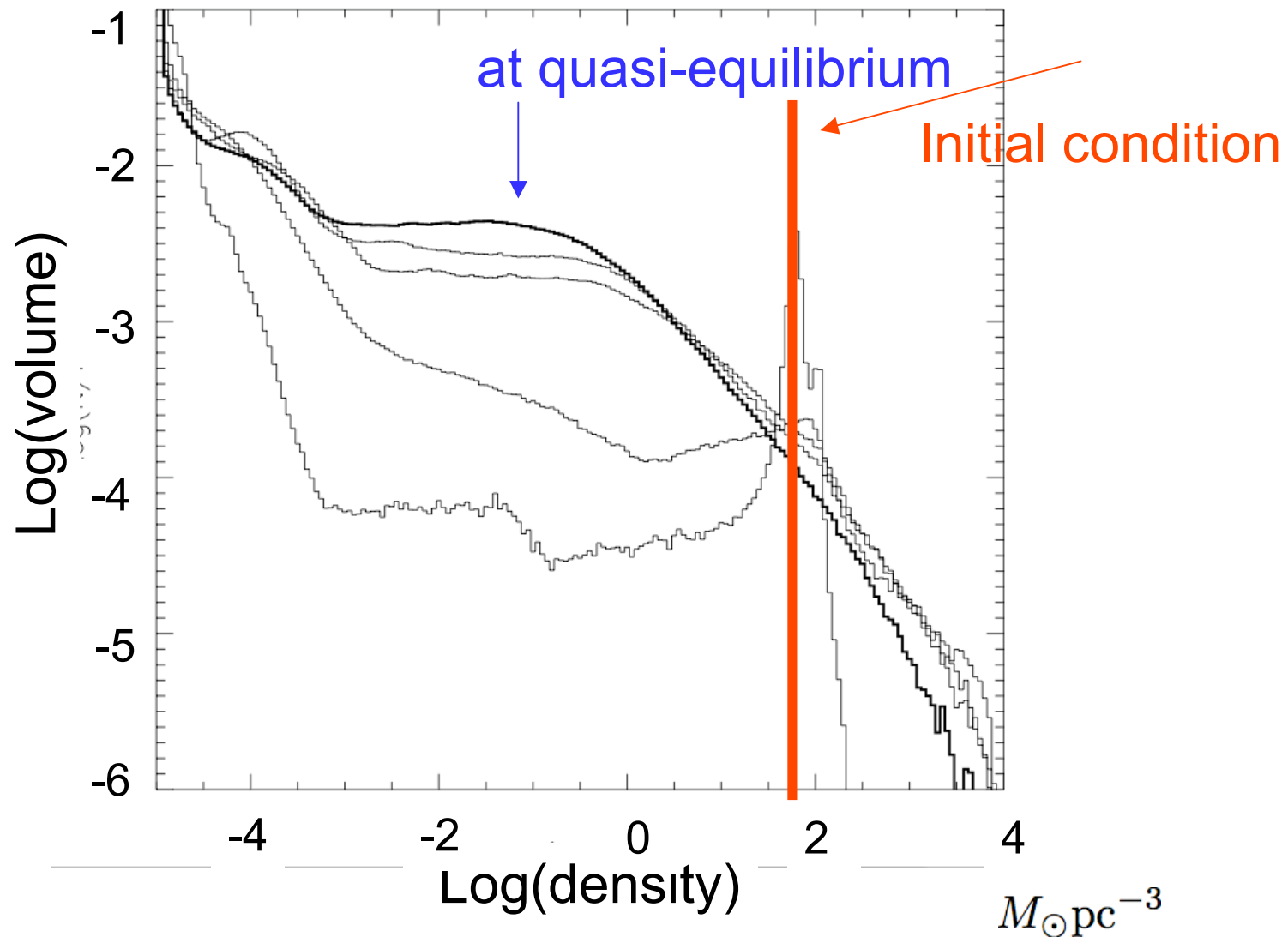


3-D evolution of a gas
disk in 2.6x2.6 kpc

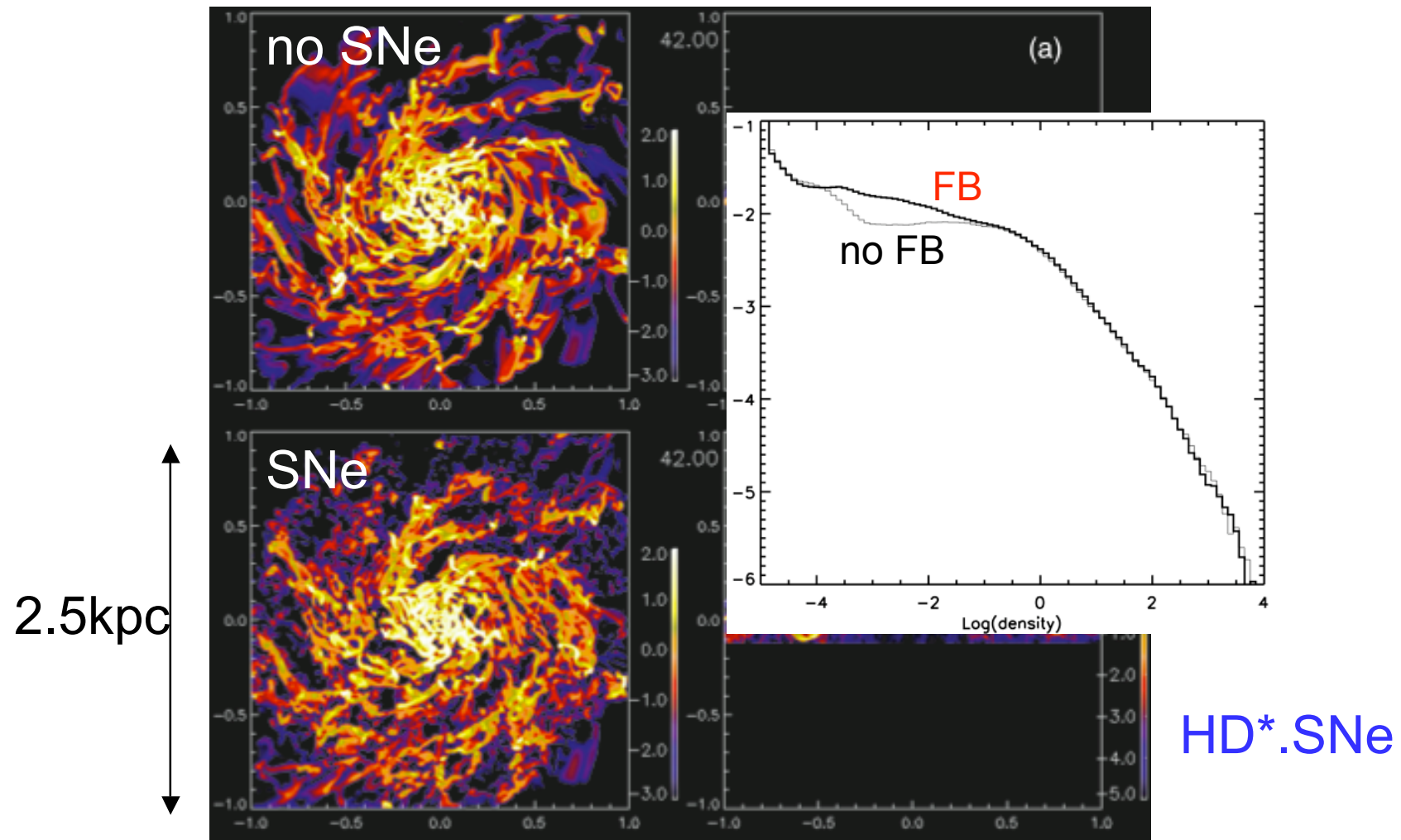
Low resolution run (10 pc)

Evolution of density PDF in the unstable disk

High-density tail and low density part are coupled, and the final PDF shows a smooth distribution.



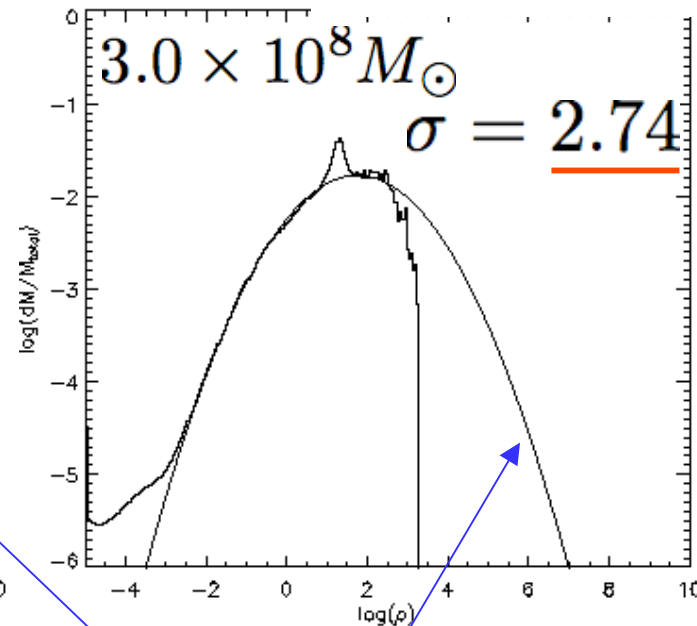
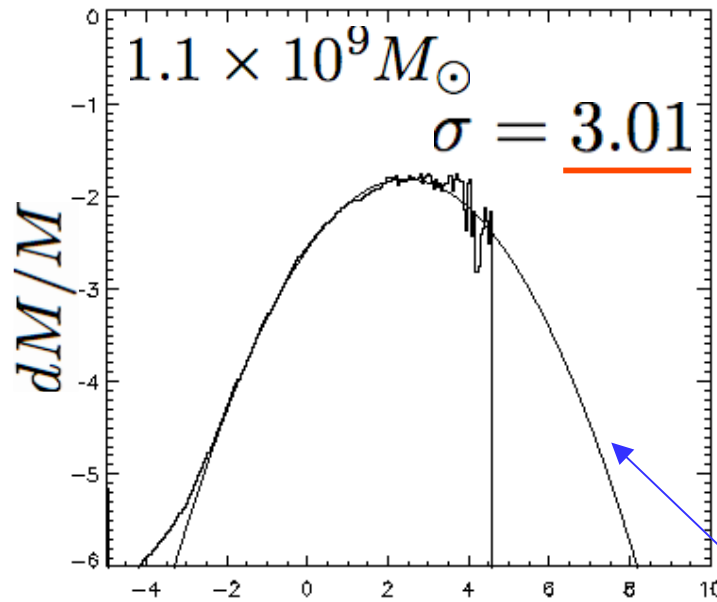
The density PDF is robust for energy feedback from SNe



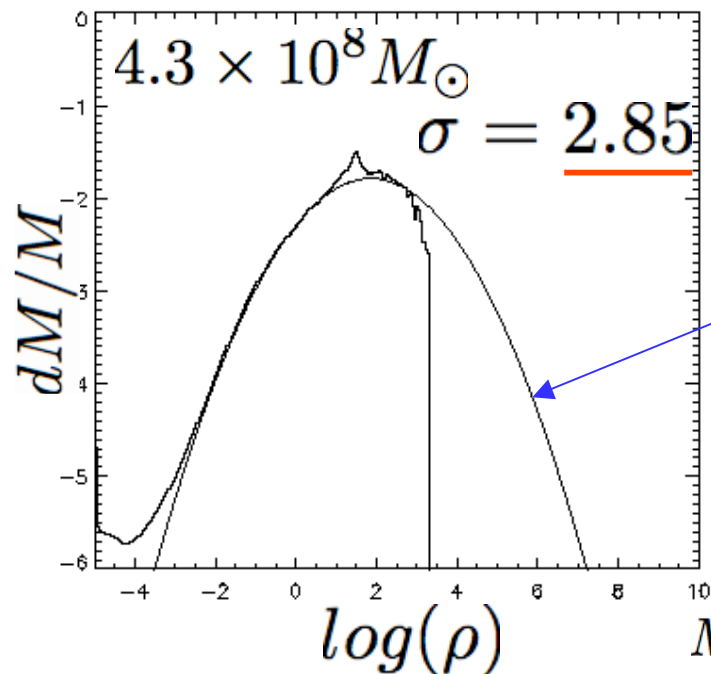
Positions of energy input are randomly selected.
SN rate= $1.5 \times 10^{-5} \text{ yr}^{-1} \text{ kpc}^{-2}$

Mass-weighted density PDF

$$f(\rho)d\rho = \frac{1}{\sqrt{2\pi}\sigma} \exp\left[-\frac{\ln(\rho/\rho_0)^2}{2\sigma^2}\right] d\ln\rho$$



In more massive disks, the LN-dispersion is larger.

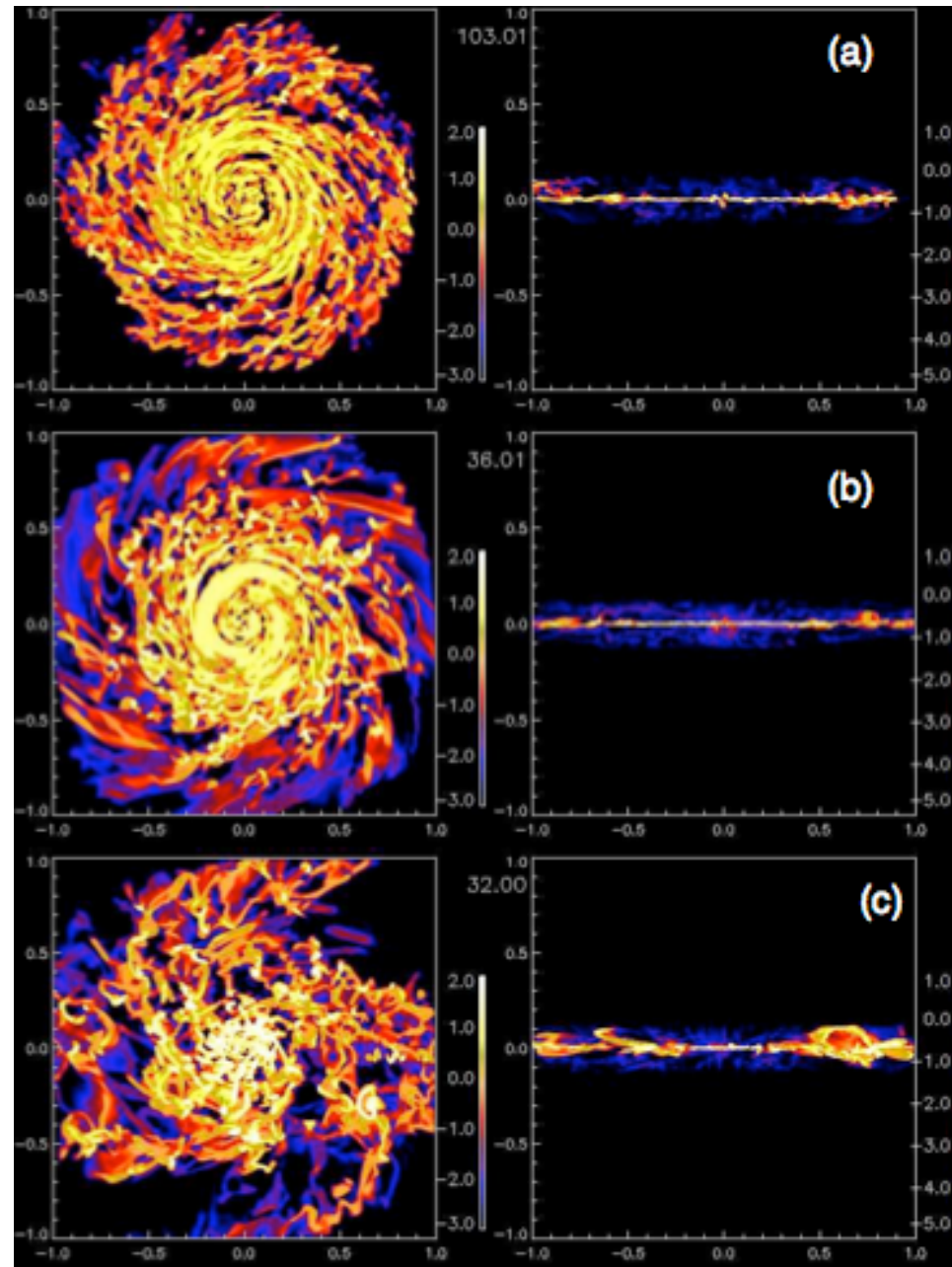
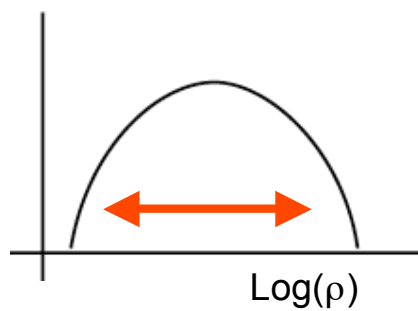
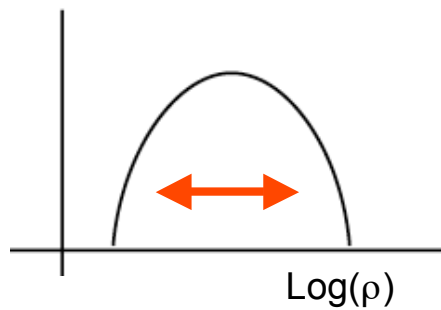
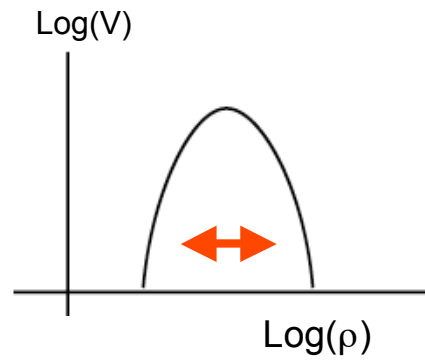


For LN-PDF

$$\sigma^2 = \frac{1}{2} \ln(\langle \rho \rangle_M / \rho_0)$$

ρ_0 is nearly constant.

$$\rho_{0,M} = \rho_{0,V} e^{\sigma^2}$$



$$M_{\text{gas}} = 3.0 \times 10^8 M_{\odot}$$

Less massive

$$4.3 \times 10^8 M_{\odot}$$

$$1.1 \times 10^9 M_{\odot}$$

massive

$$\langle \rho \rangle_V = \rho_0 e^{\sigma^2/2}$$

Larger dispersion in more massive system

star formation model based on LN-PDF

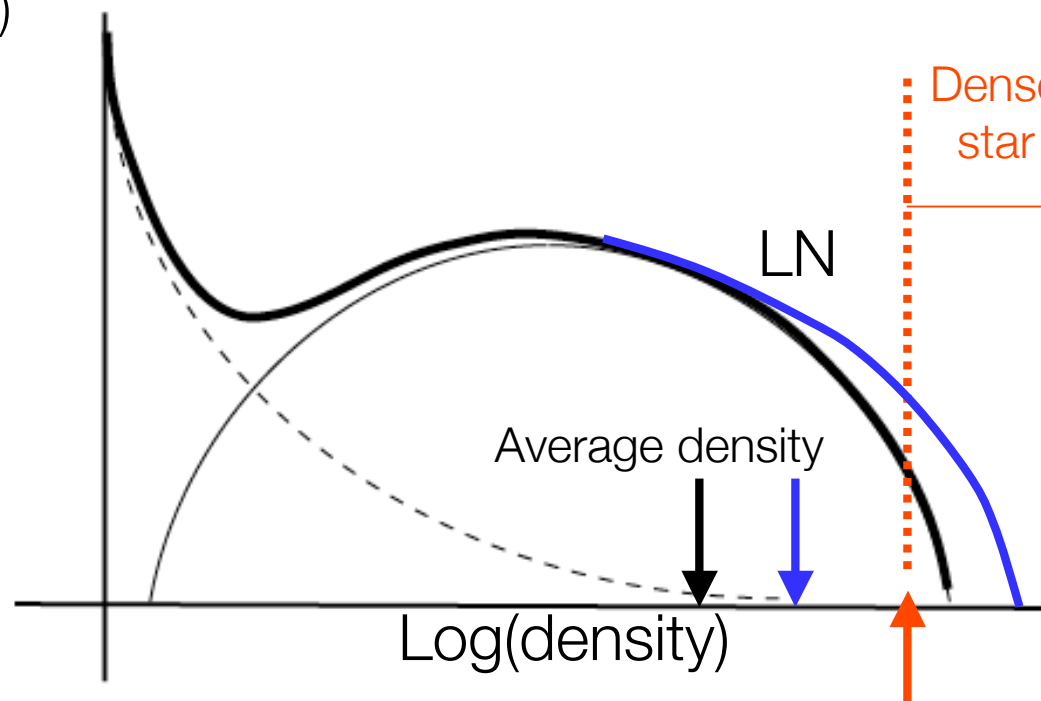
Elmegreen (2002), KW & Norman (2007)

Average density $\uparrow \Leftrightarrow$ dispersion of LN \uparrow

cf. Krumholz & McKee 05 for
molecular clouds

\Rightarrow fraction of high density clumps $\uparrow \Rightarrow$ SFR \uparrow

Log(Volume)



threshold density for
local SF

e.g. $n=10^{3-5} \text{ cm}^{-3}$

$$\langle \rho \rangle_v = \rho_0 e^{\sigma^2/2}$$

SFR is scaled to gas density averaged on a kpc-scale

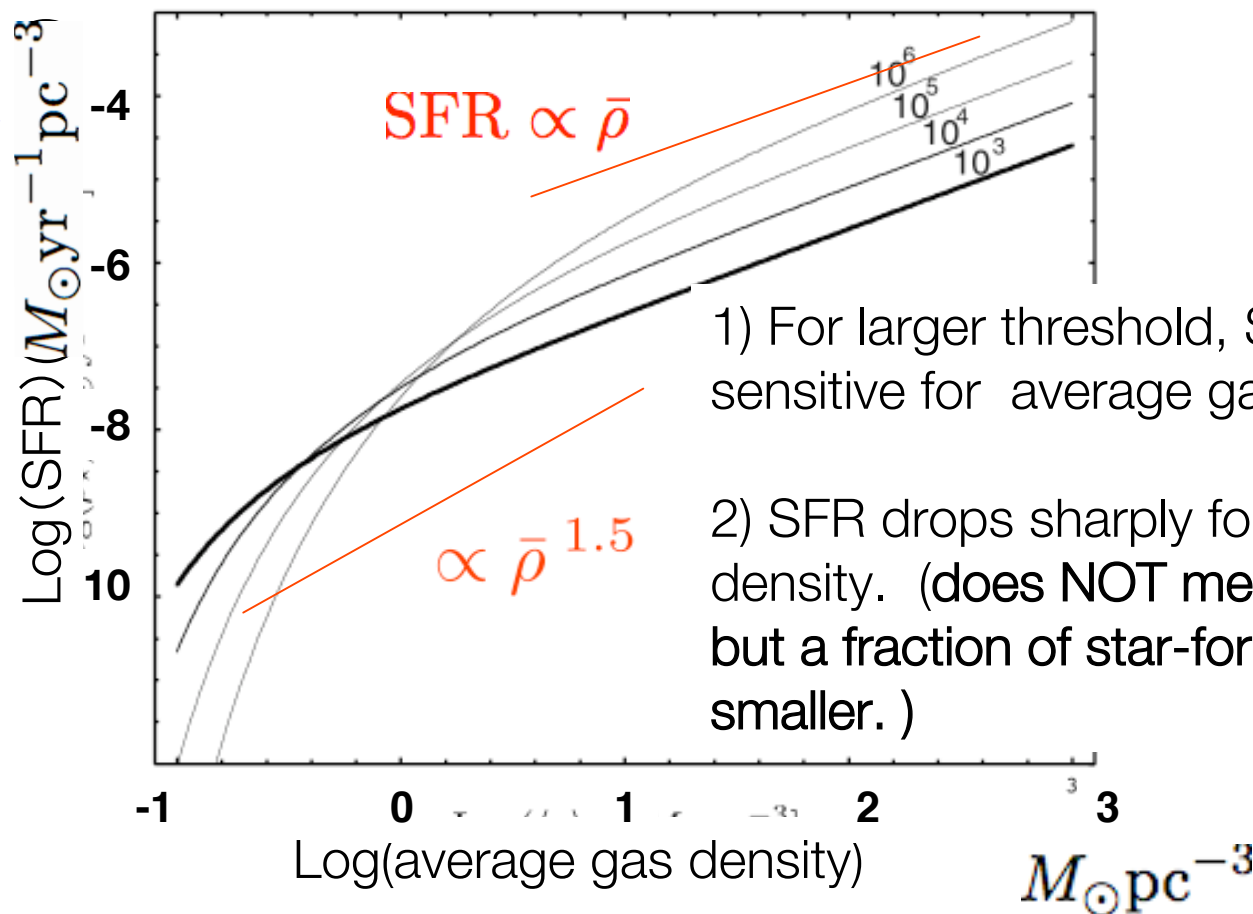
KW&Norman07

If stars are formed in a free fall time above
the threshold density,

$$\dot{\rho}_\star = \epsilon_c (G\delta_c)^{1/2} f_c \rho_0^{3/2} e^{\sigma^2/2}$$

Threshold density = $10^6, 10^5, 10^4, 10^3 \text{ cm}^{-3}$

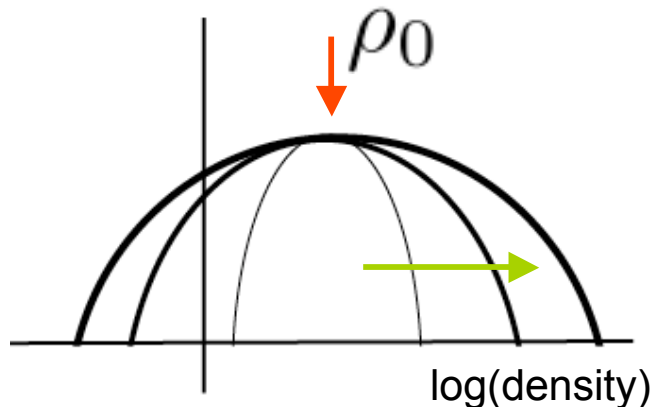
efficiency



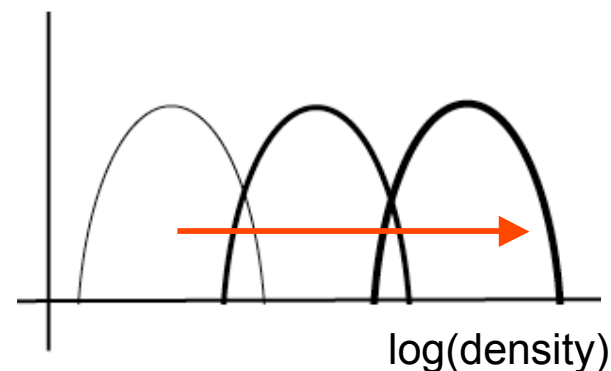
How does PDF change depending on total gas mass (average density)?

- Suppose the density distribution is represented by LN-PDF, the PDF responds to increasing the total gas mass by two ways:

(1) Characteristic density is const., & dispersion increases.



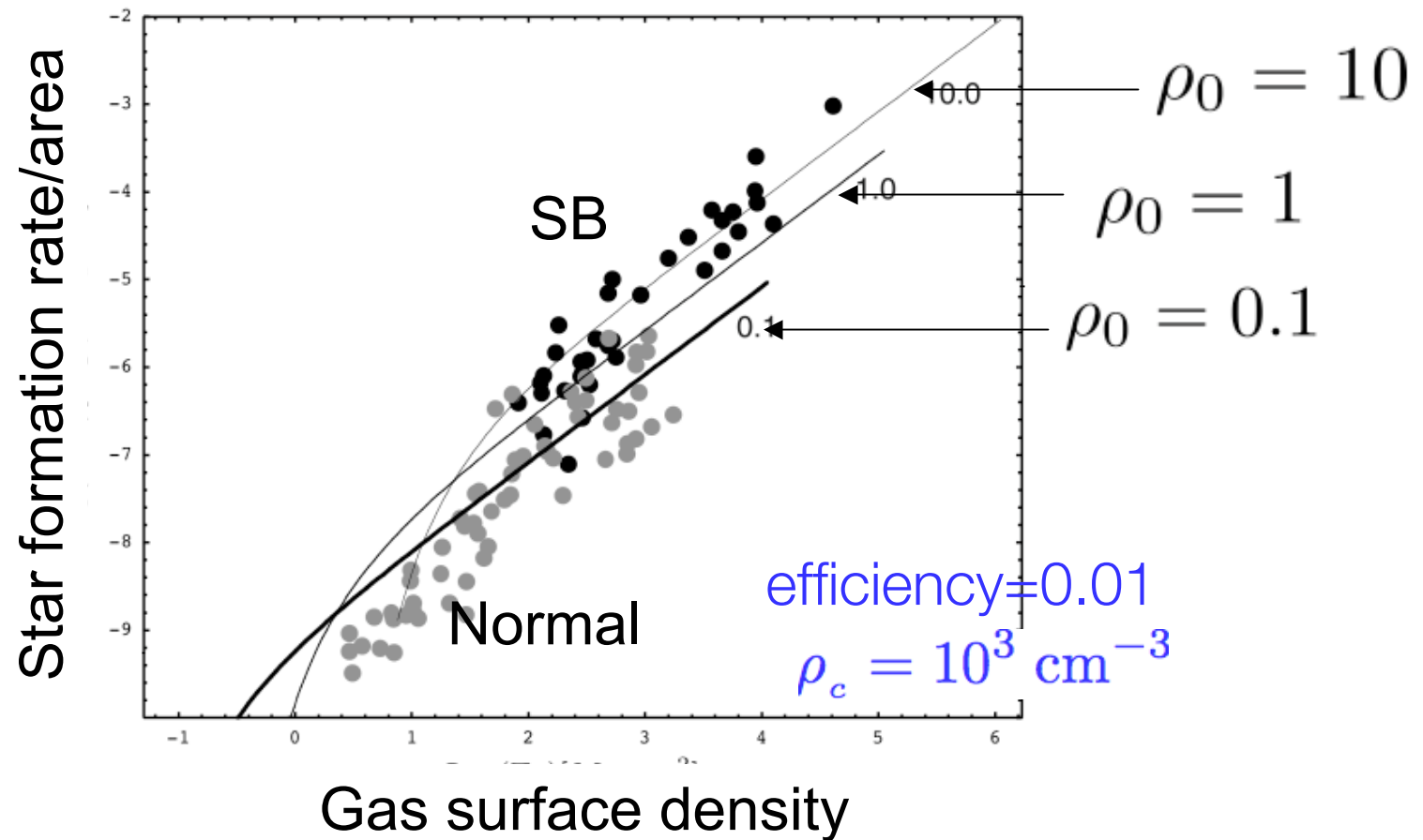
(2) Dispersion is constant, & Characteristic density increases.



Parameter to fit observed SFR: ρ_0 ϵ_c ρ_c

Comparison with observed SFR

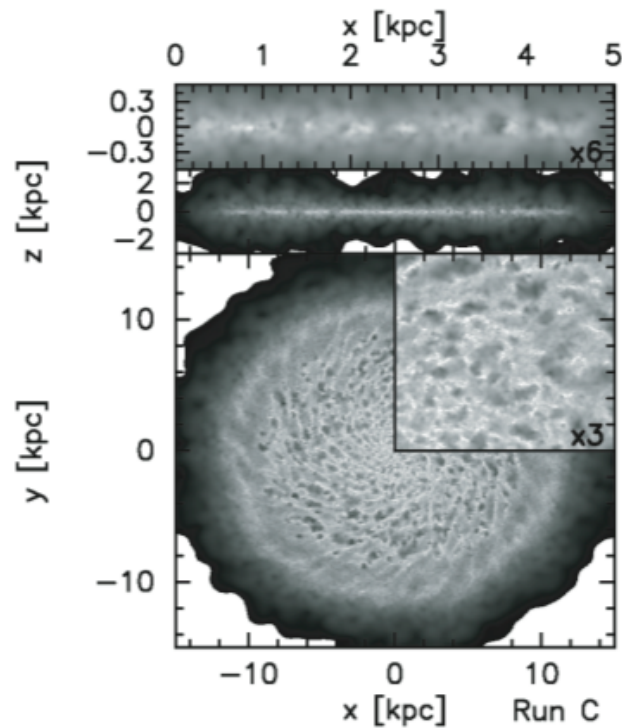
In starburst galaxies, the characteristic density (ρ_0) should be larger, for a given efficiency and threshold density of local star formation.



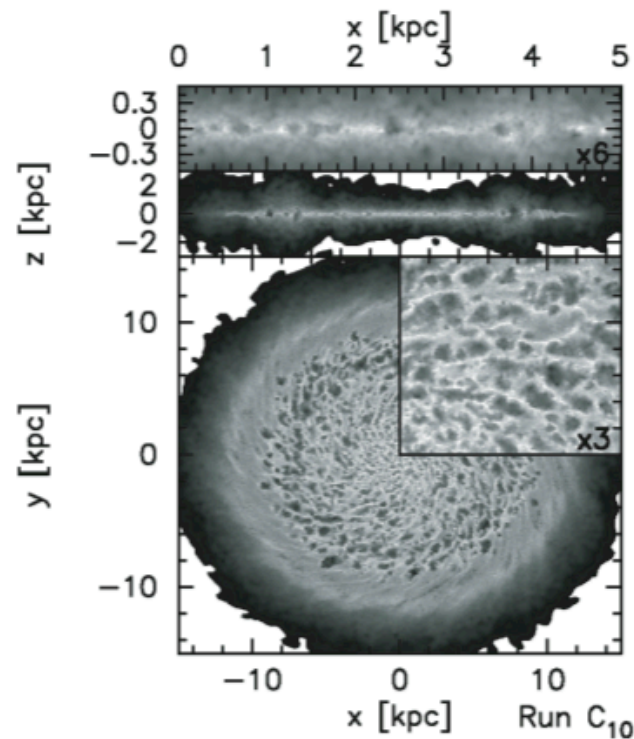
Observed data: Komugi + (2005) based on CO survey in Nobeyama

N-body/SPH simulations (*ASURA*)

Saitoh et al. (2008) PASJ 60, 667

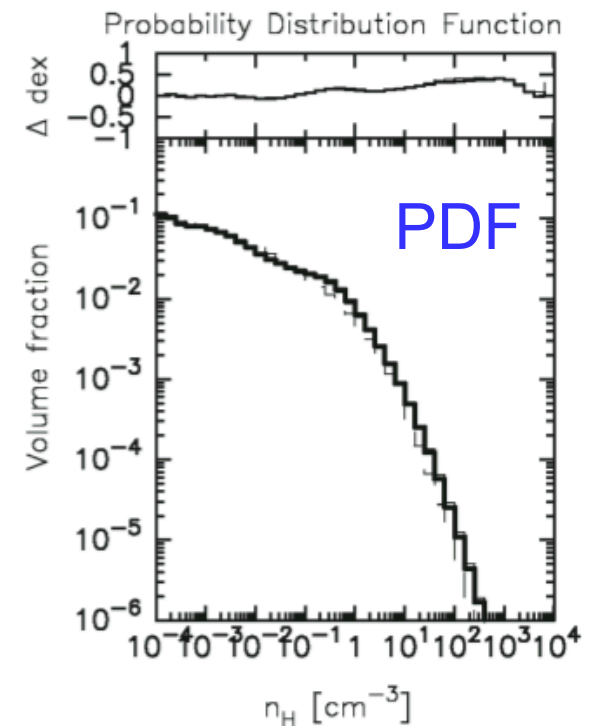


$$N_{\text{SPH}}=10^6$$

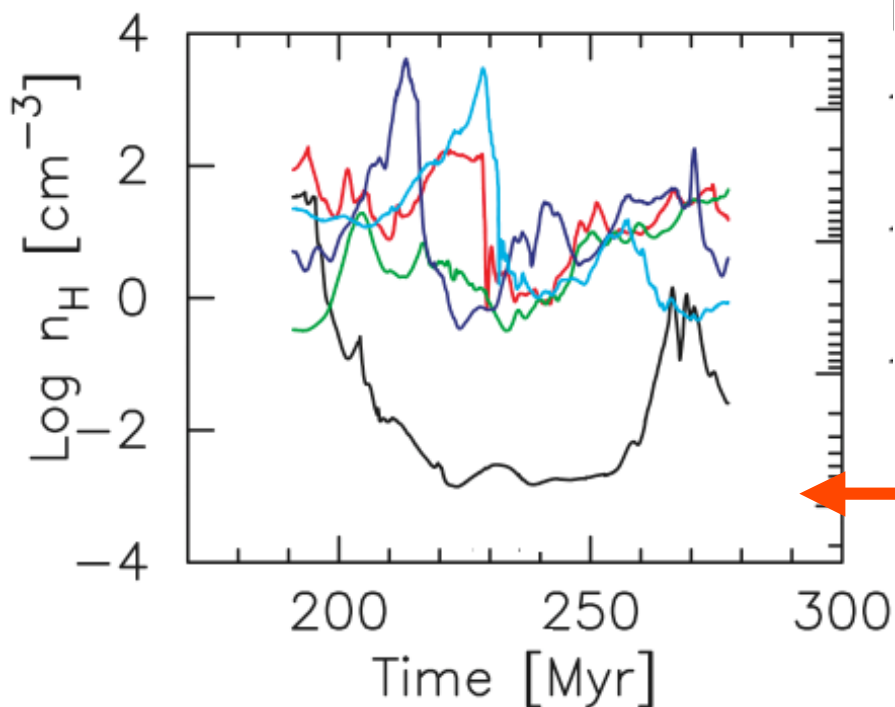
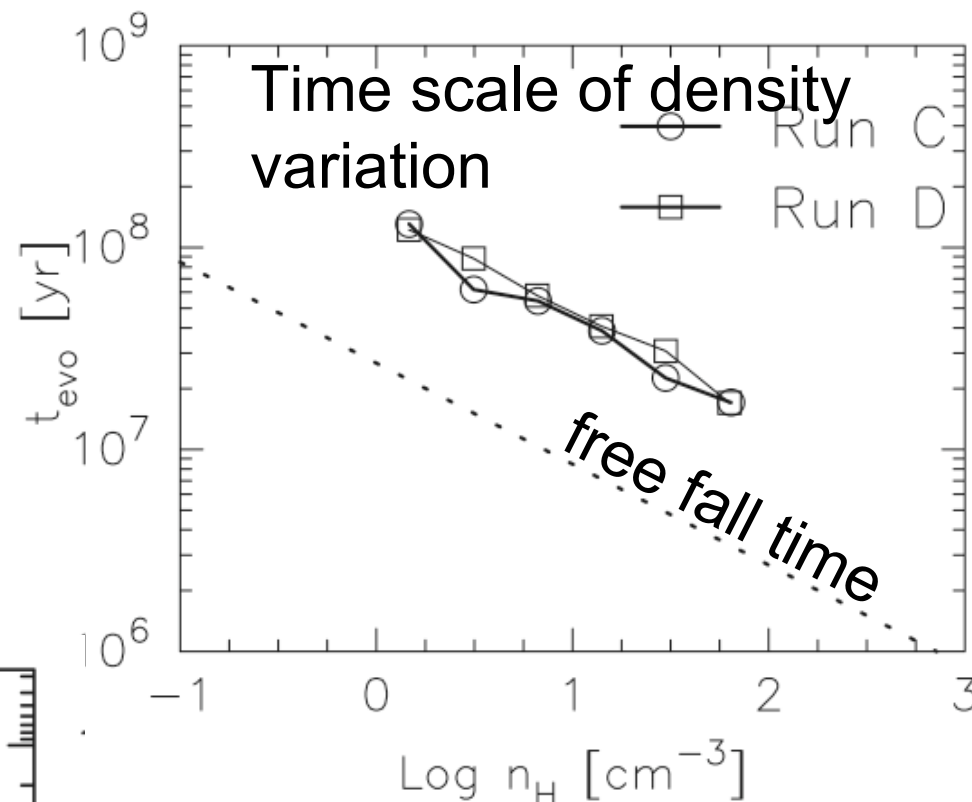
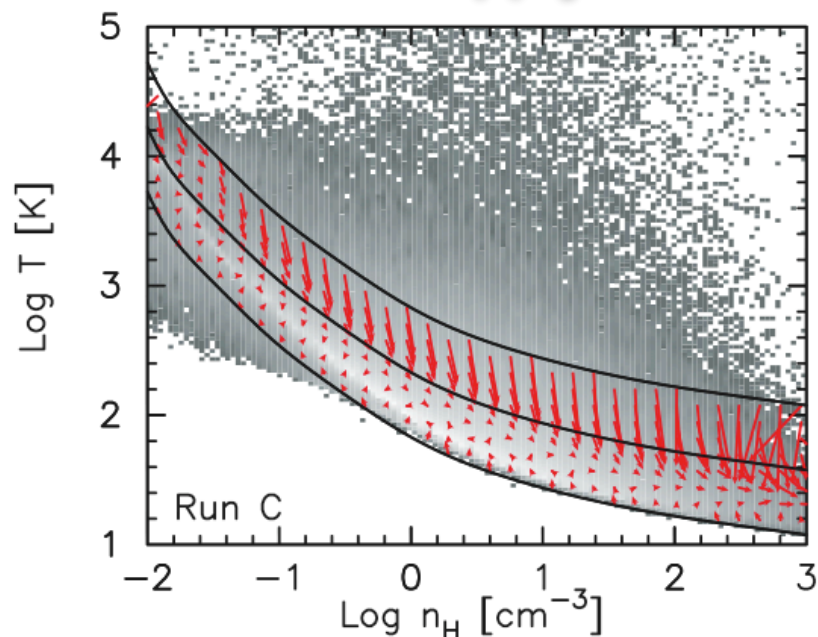


$$N_{\text{SPH}}=10^7$$

Density structures converge



Mass supply toward high density region is slow!



Net evolution of density is
 $\sim 5 \times$ slower than free fall time on
average

“Random walk” in density

Density and temperature distribution around a SMBH

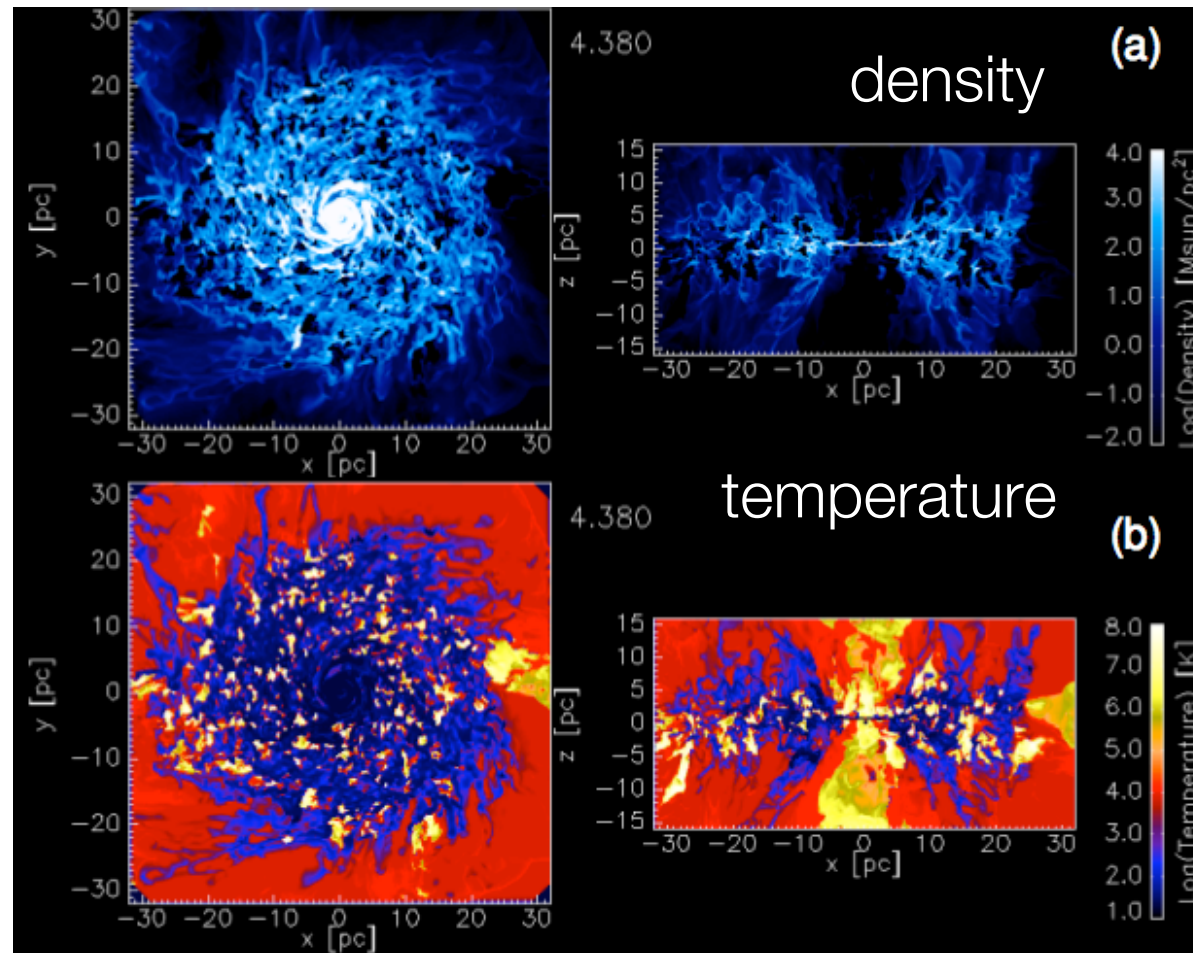
$$M_{\text{BH}} = 1.3 \times 10^7 M_{\text{sun}}$$
$$M_{\text{gas}} = 6 \times 10^6 M_{\text{sun}}$$

Wada, Papadopoulos, Spaans (2009) arXiv:0906.5444

HD*.H2.SMBH.SNe

Resolution: 0.125 pc

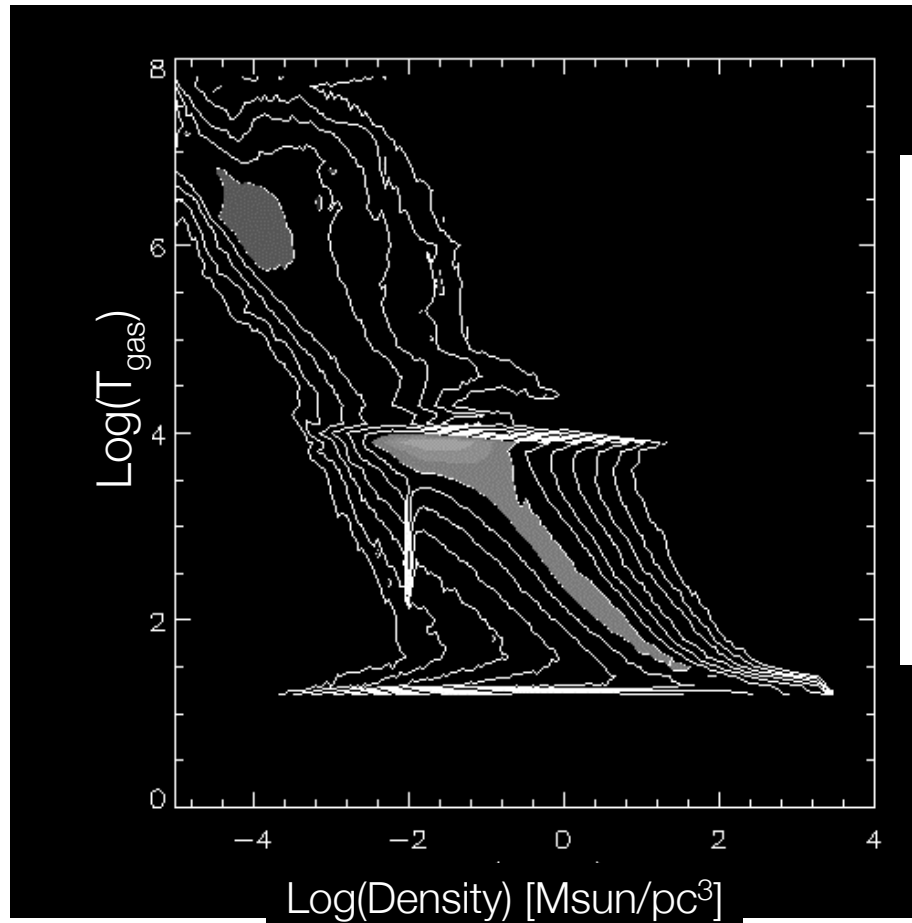
64 pc



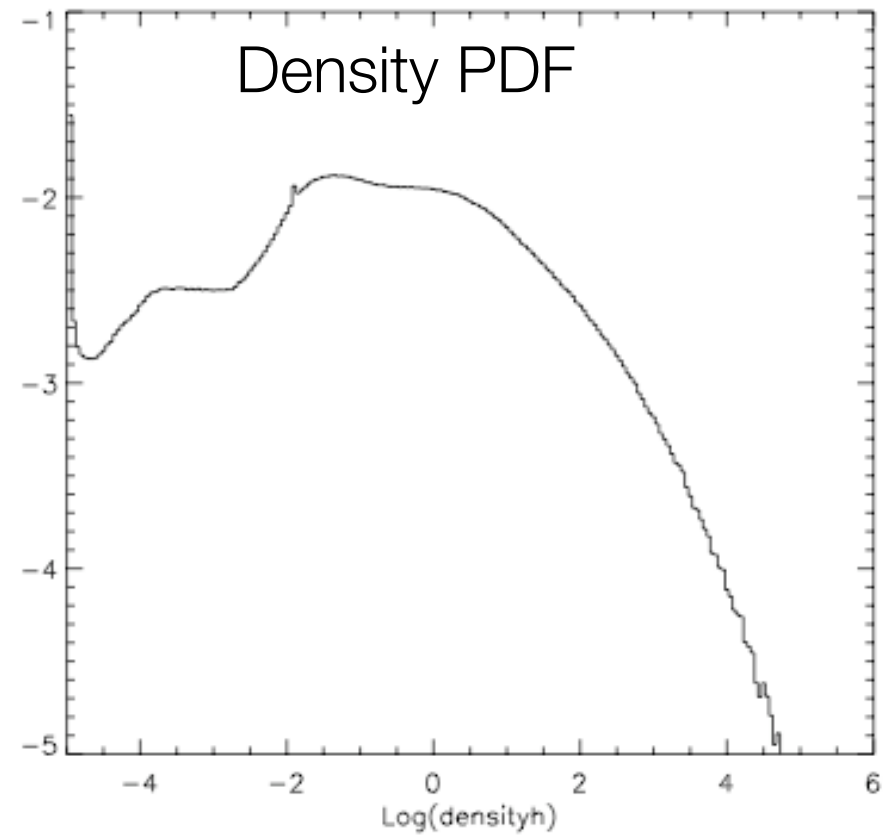
$$G_0 = 10$$

$$\text{SN rate} = 5 \times 10^{-5} / \text{yr}$$

Phase-diagram and PDF in a galactic central region



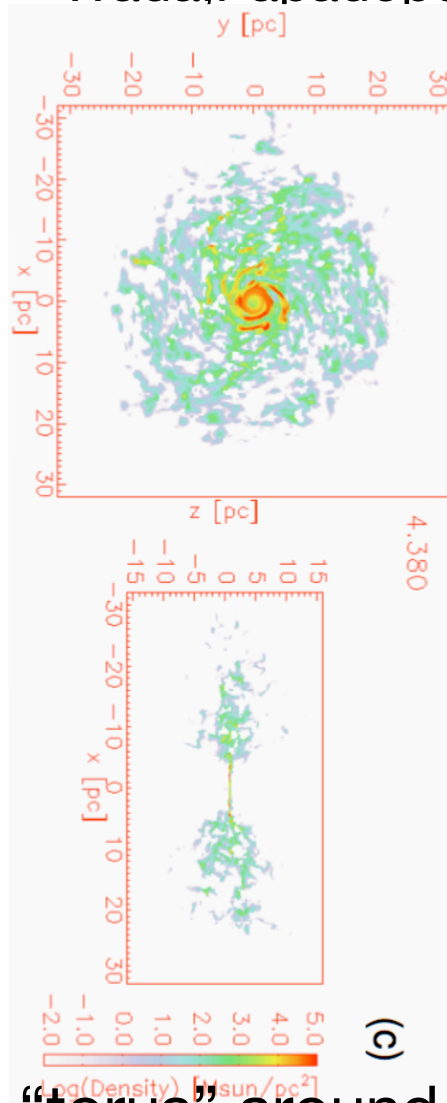
$\text{Log}(\text{Volume fraction})$



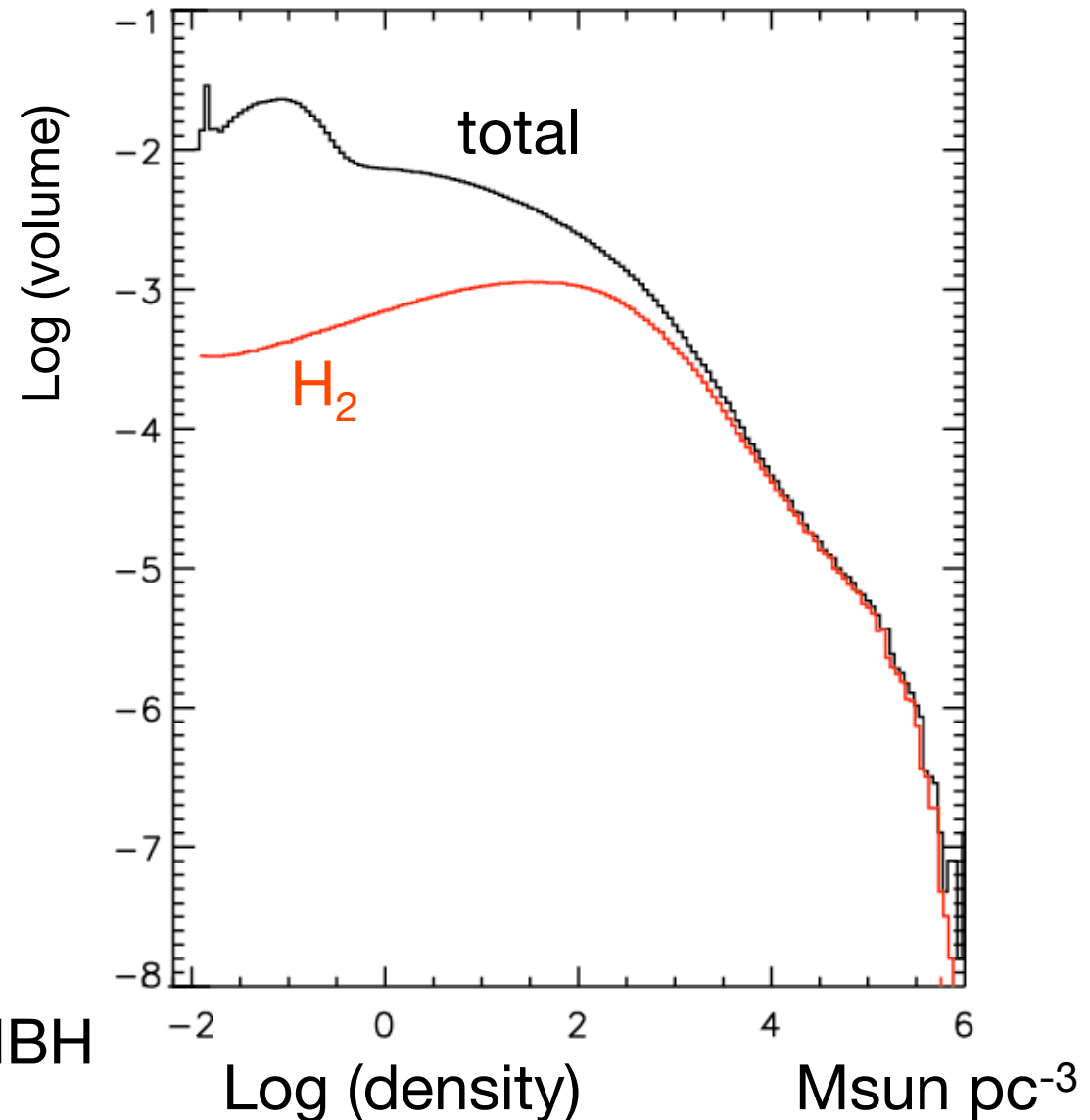
Msun pc^{-3}

Density & H_2 PDF in a galactic central region

Wada, Papadopoulos, Spaans (2009) arXiv:0906.5444



H_2 “torus” around a SMBH



summary

- The **inhomogeneous** ISM in **globally stable** disks can be characterized by a **Log-Normal like density PDF**
 - Global SFR is controlled by a mass fraction in high density regions, which is NOT independent of lower density regions.
 - Time-scale of global SF is not determined only by the free fall time of high density gas, but also by evolutionary time scale toward high density regions.
 - Mass ‘flow’ is slow in a statistical sense due to diffusion-like evolution in a density space.
- Non-LN PDF or multi-component PDF would be also the case in realistic situations.

